



Wright  
Laboratory

# Jet substructure and hadronization with STAR

Raghav Kunnawalkam Elayavalli (Yale)  
For the STAR Collaboration

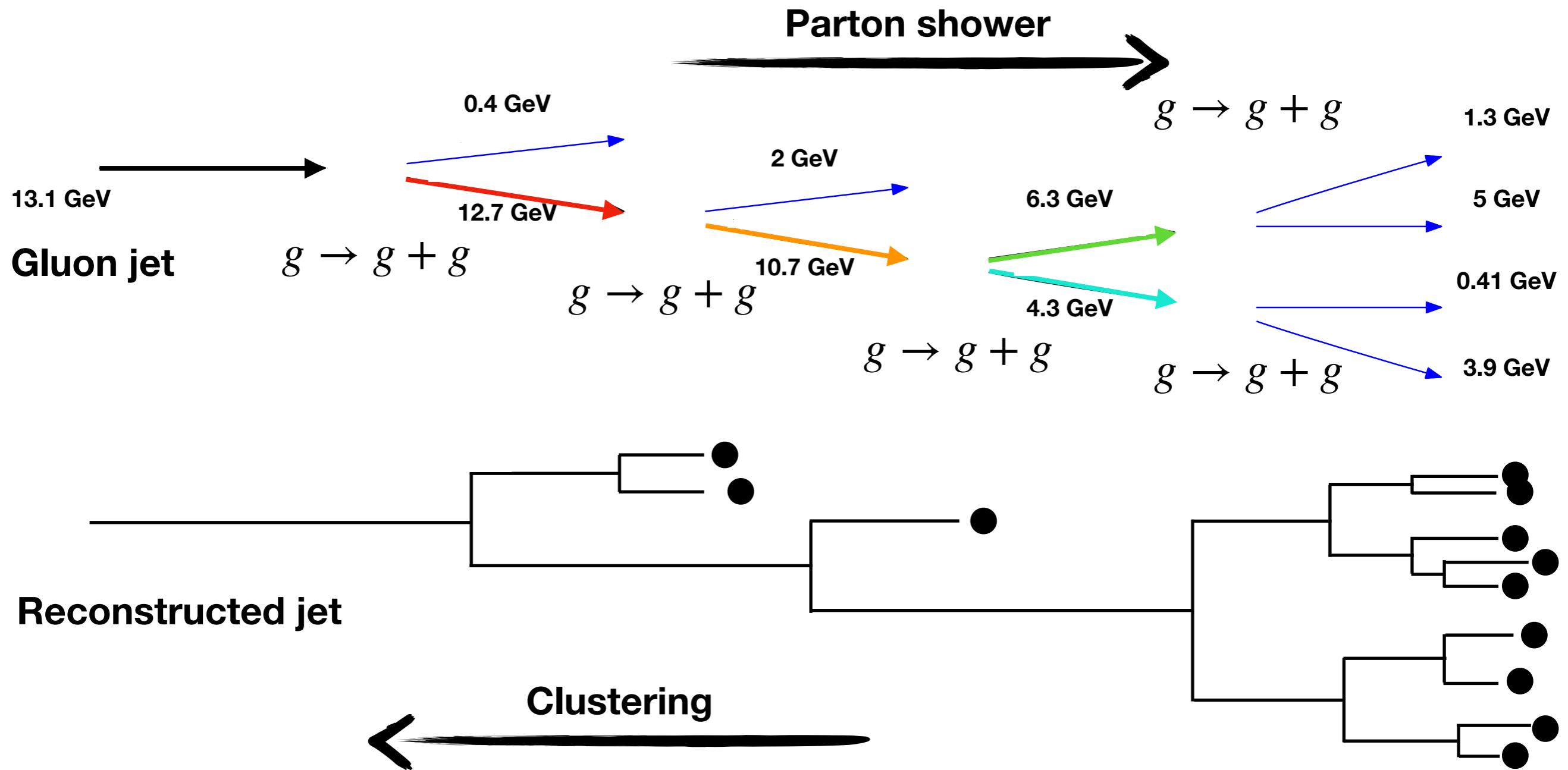
2<sup>nd</sup> Workshop on Jets for 3D Imaging at the EIC



Sept 27-29, 2021

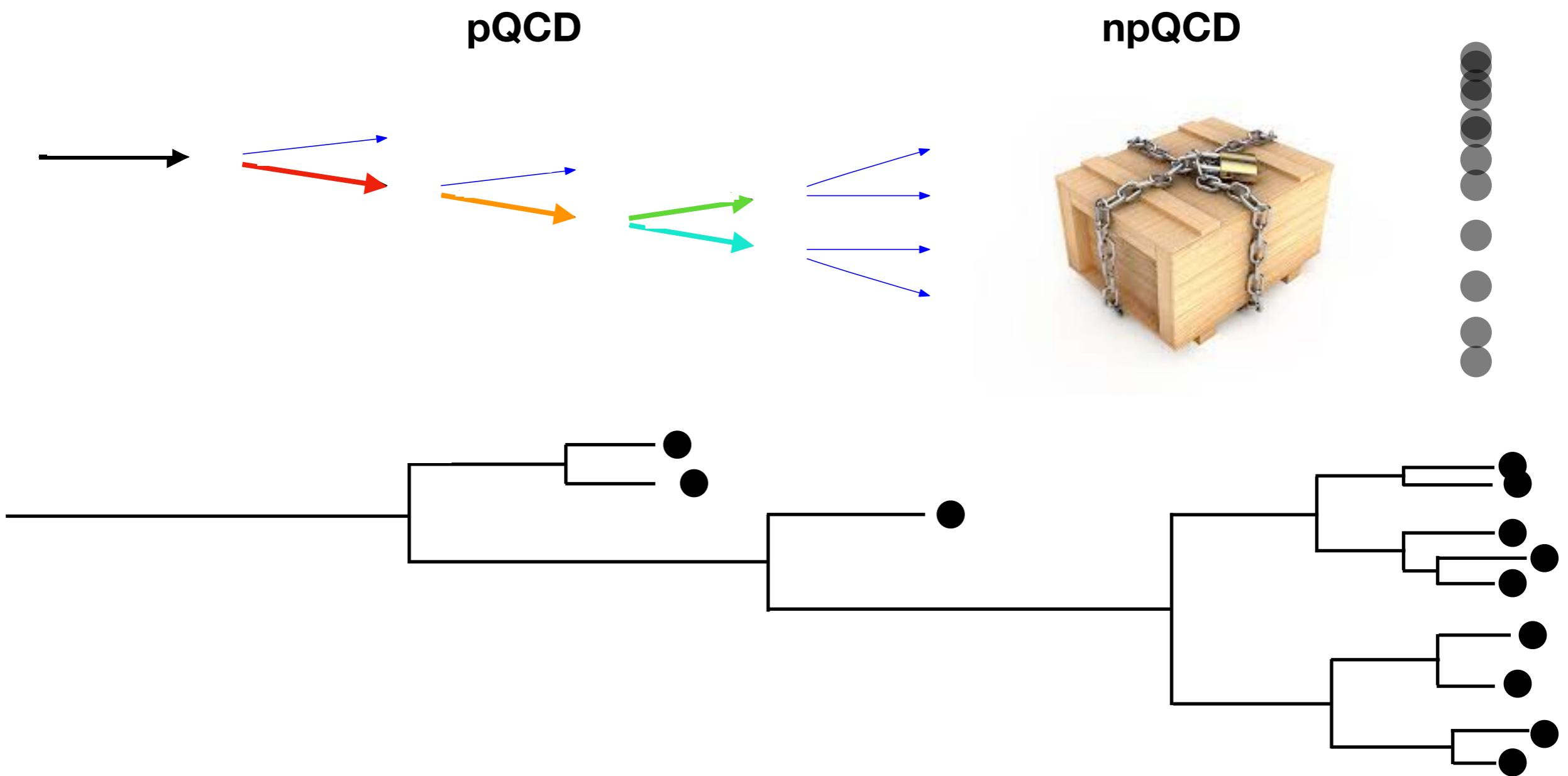
# What do we want to measure?

We want to translate an intrinsic (and unmeasurable) parton shower to experimentally accessible observable(s)

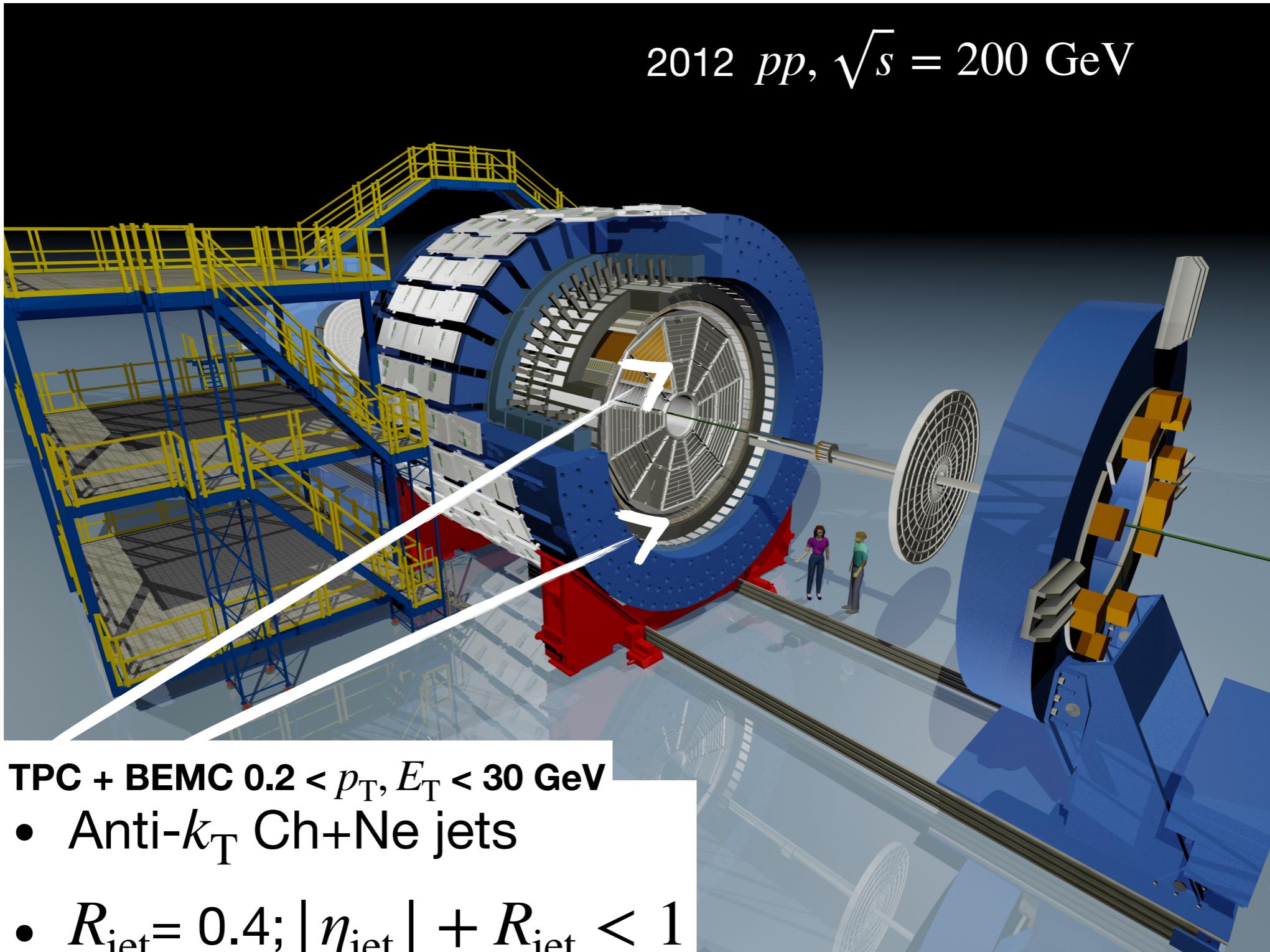


# What do we want to measure?

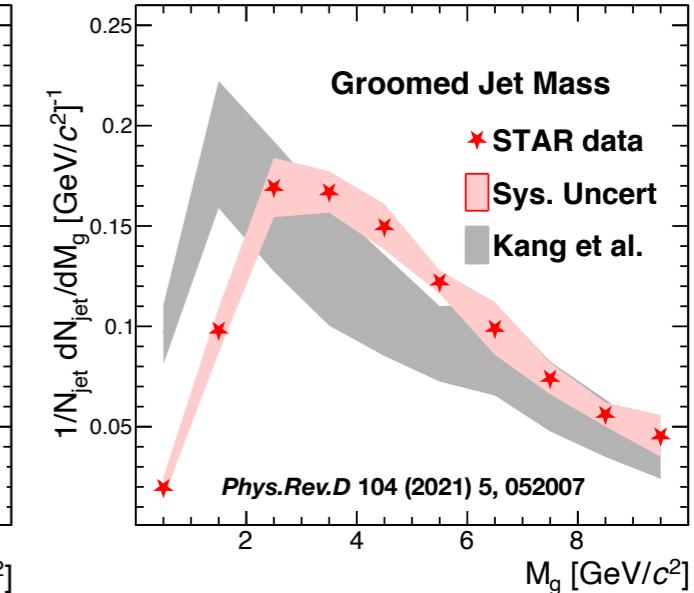
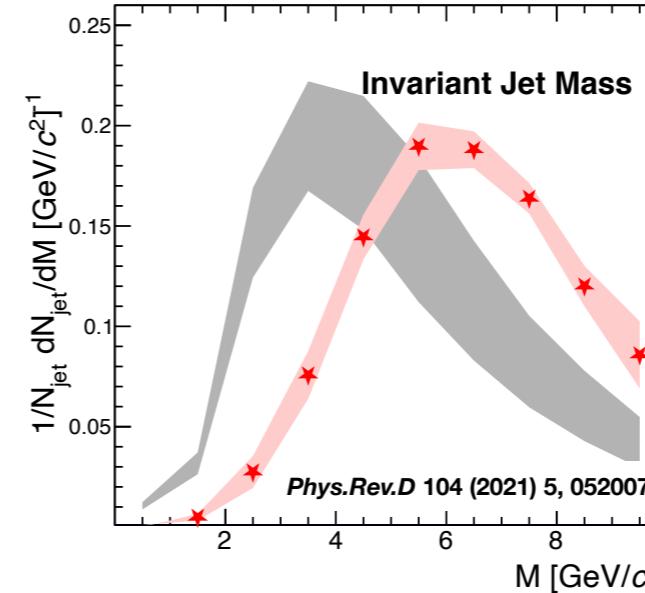
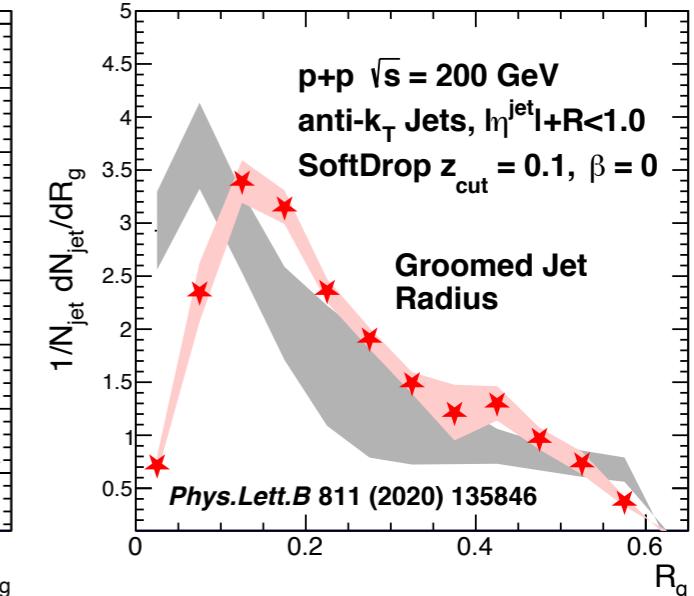
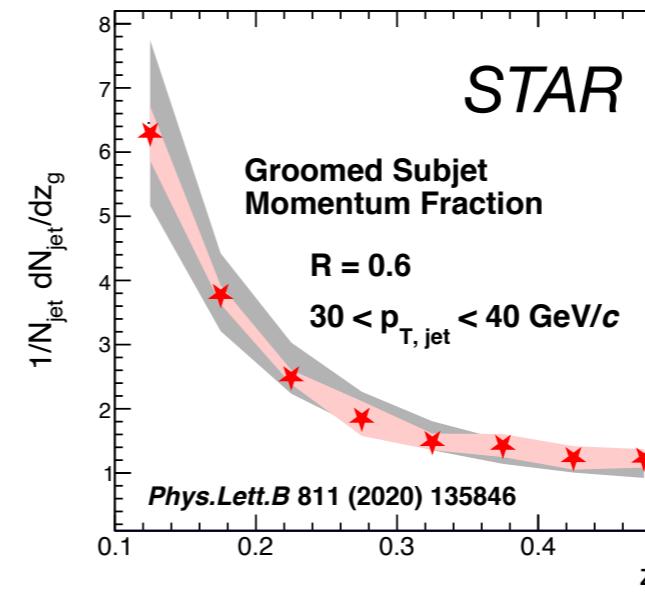
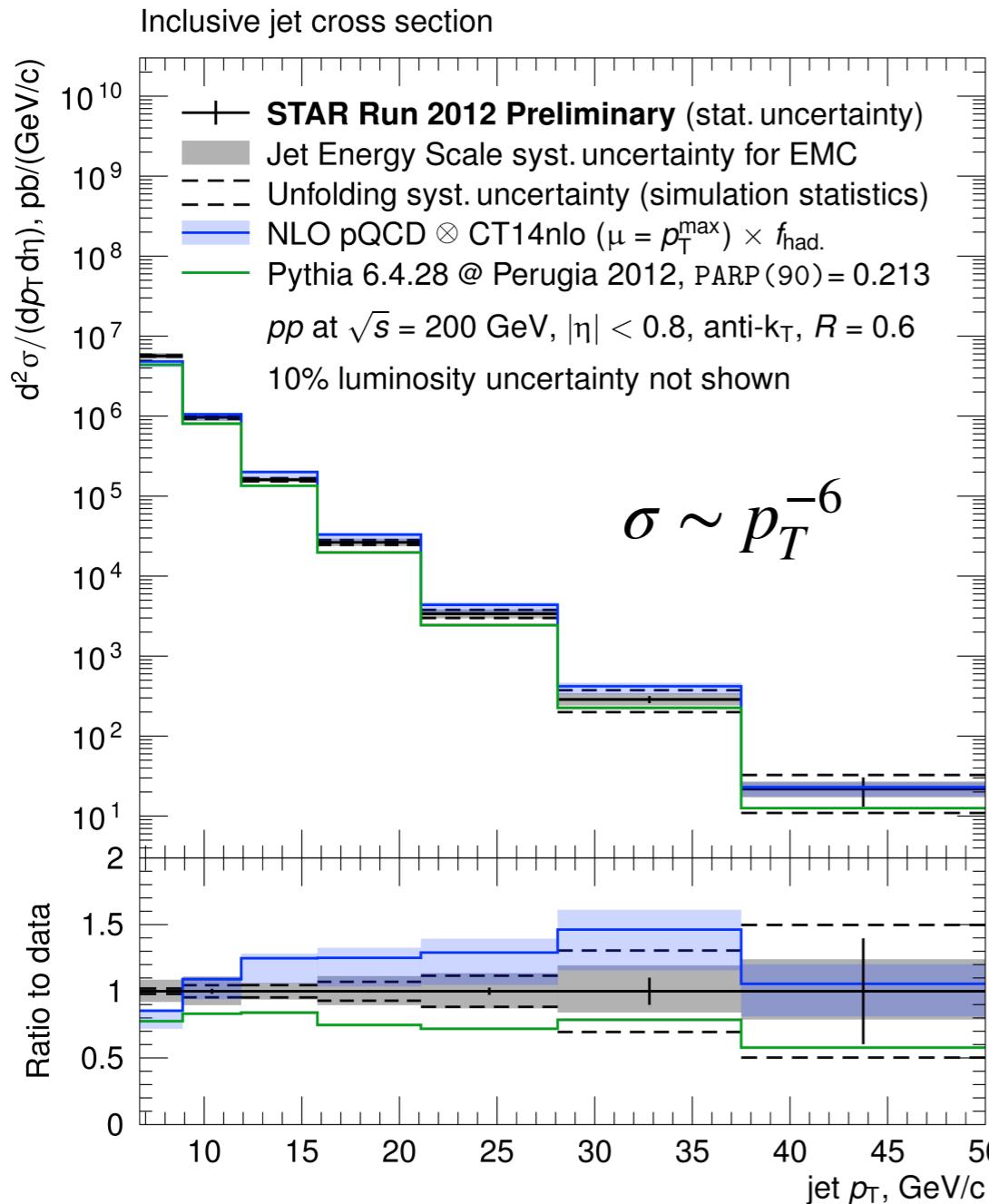
Isolate regions of phase-space within jet evolution that are sensitive to pQCD vs npQCD



# Jet reconstruction at STAR



# Jets in $pp$ $\sqrt{s} = 200$ GeV

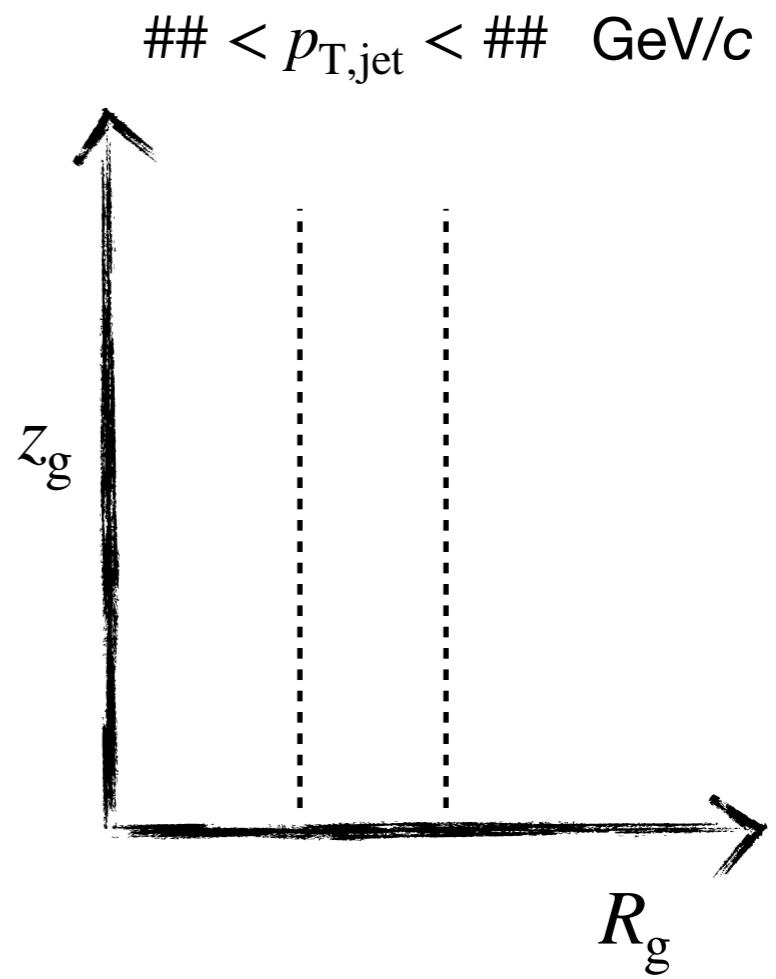
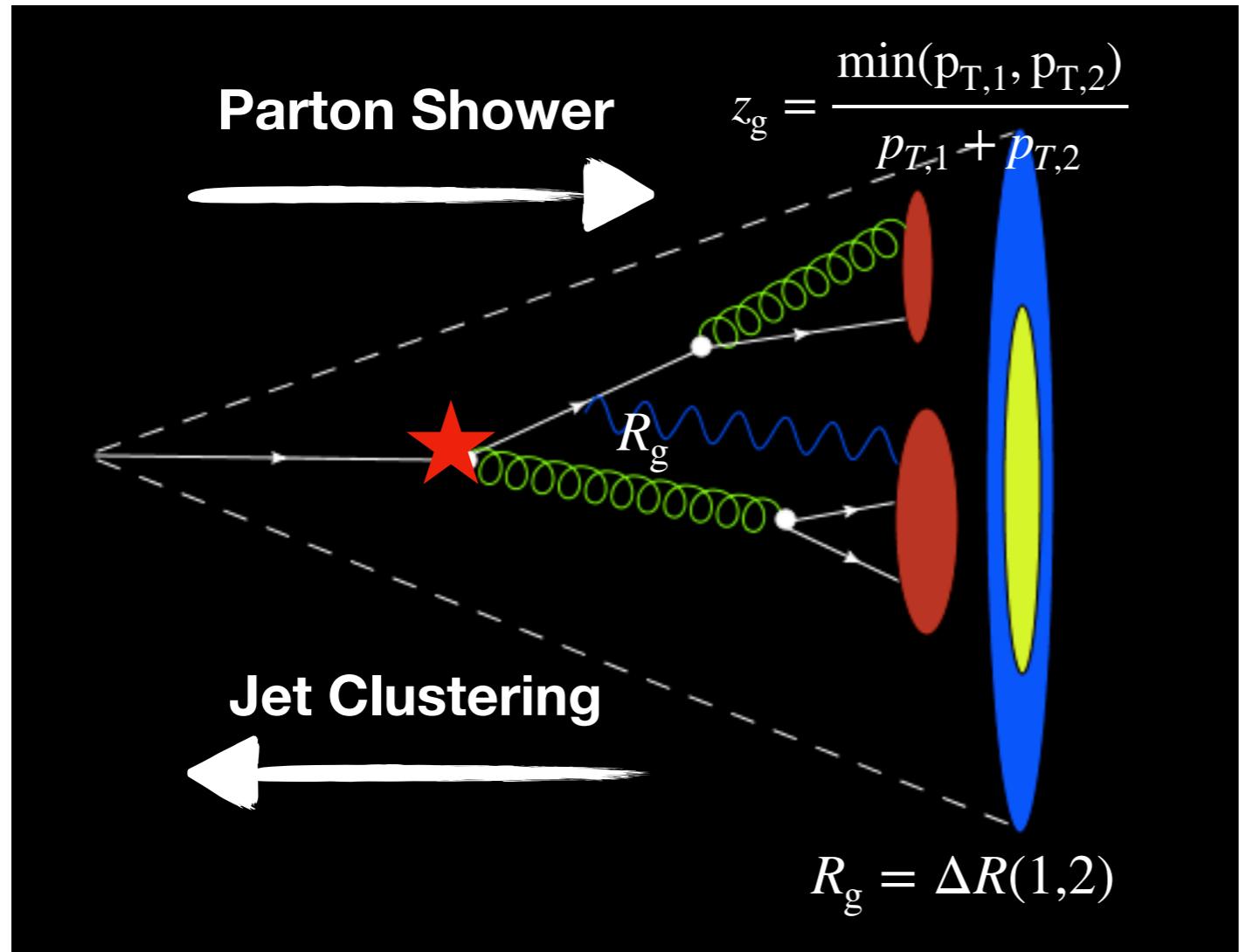


Unique population of jets with varied substructure  
Scales extend from jet  $p_T \rightarrow \Lambda_{\text{QCD}}$  (similar to the EIC)

Correlations between  
substructure observables  
at the first split

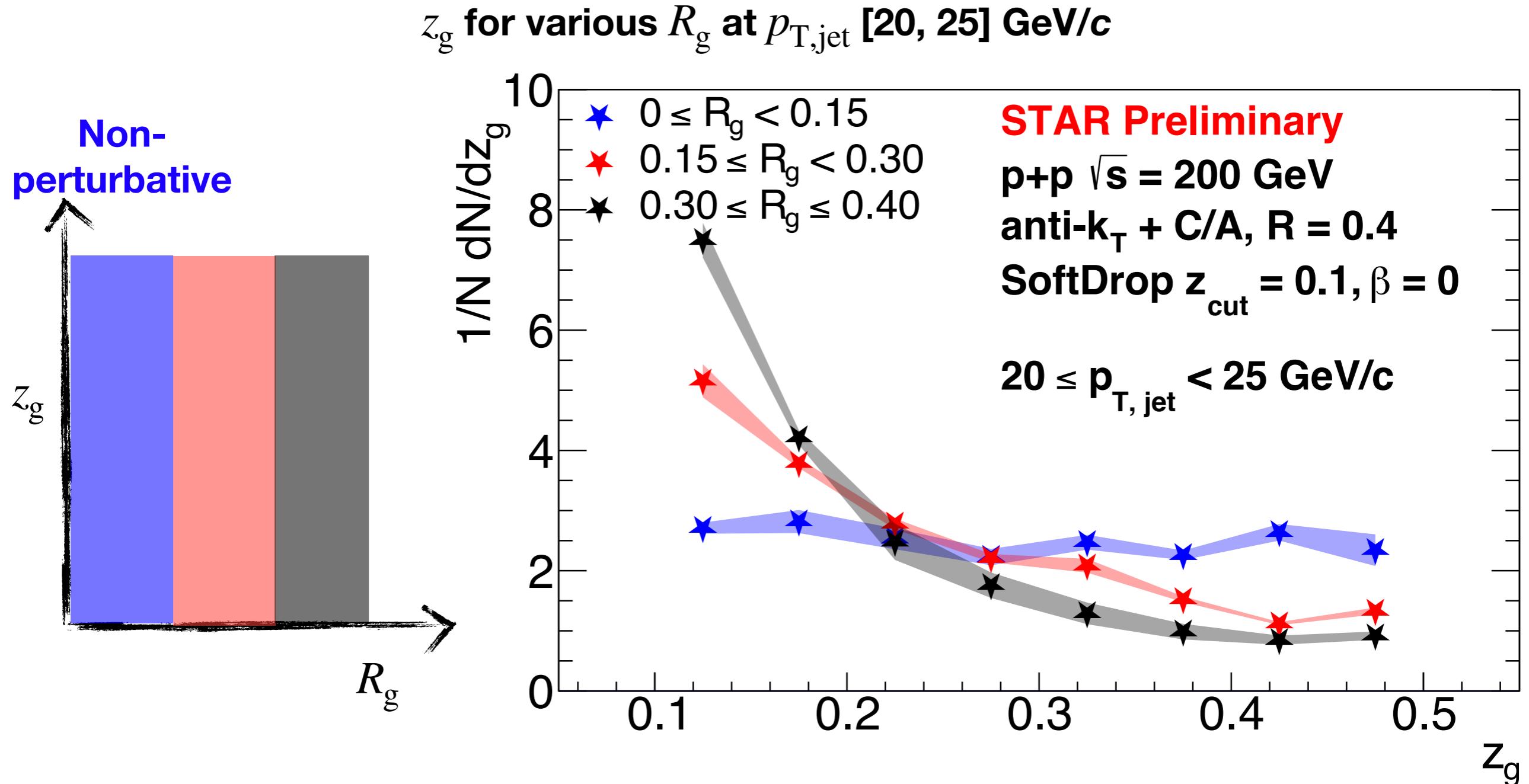
Evolution of the splitting  
observables as we travel  
along the jet shower

Formation time connecting  
jet splitting and charged  
particles in the jet



Correlations between  
substructure observables  
at the first split

# Fully corrected results

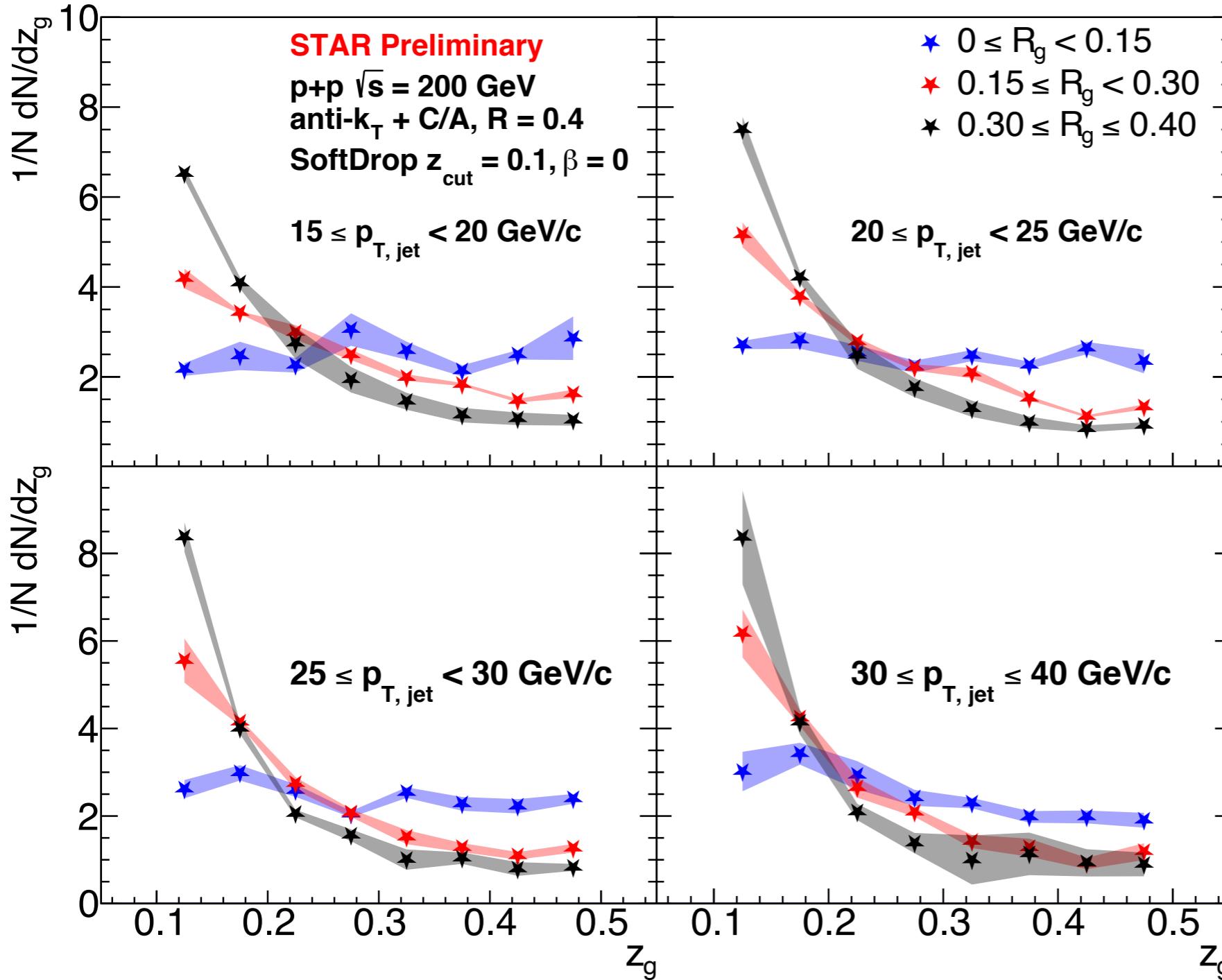


- Significant dependence on R<sub>g</sub>
- Evolution from **soft-wide angle splits** to **hard-collinear splits**

# Evolution VS.

$p_{T,jet}$

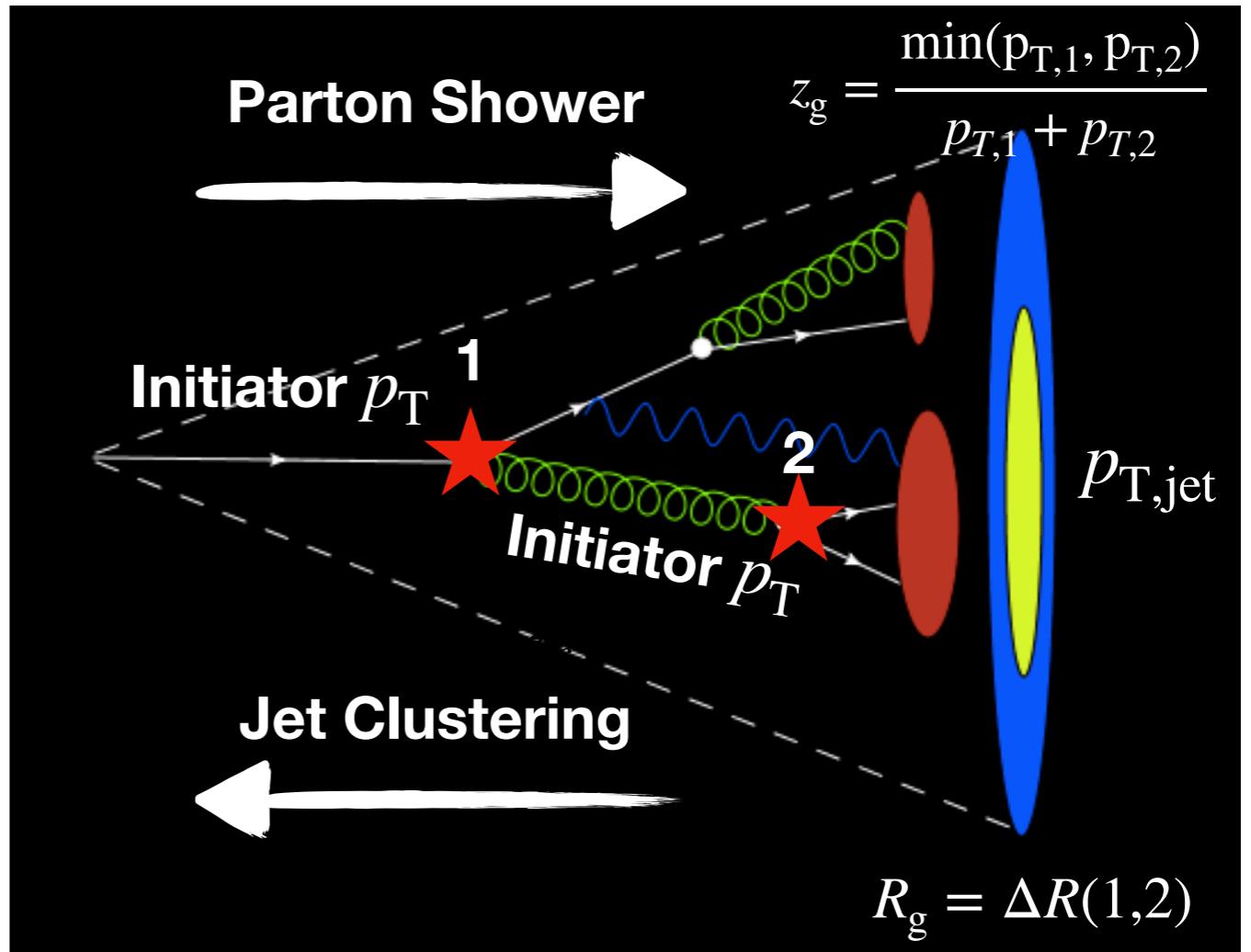
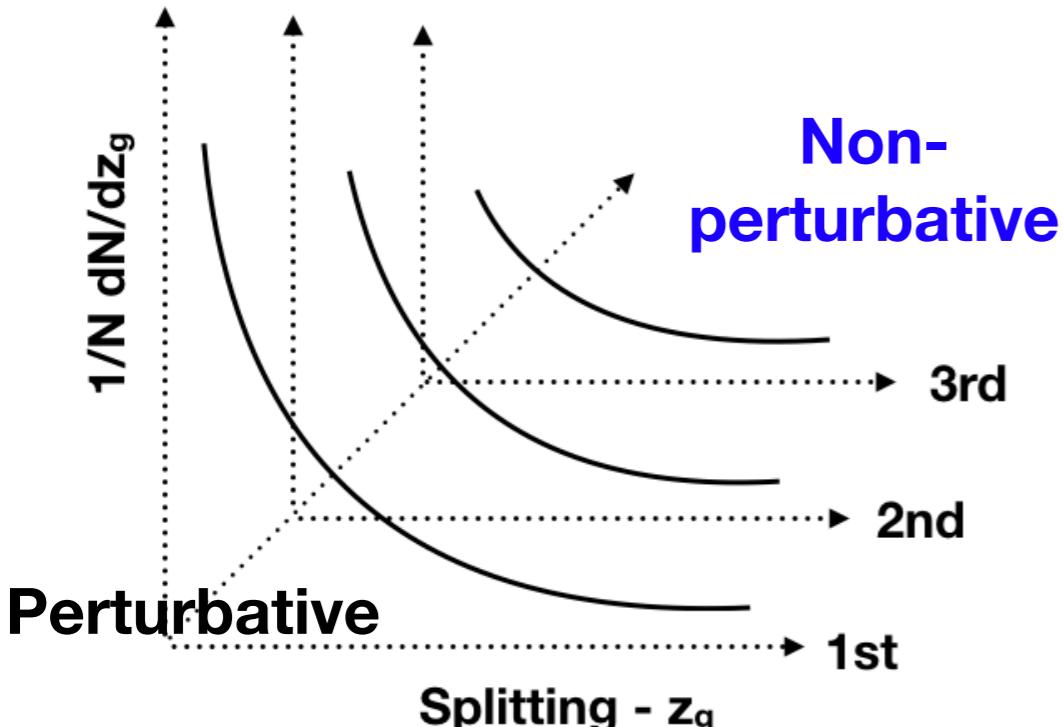
$z_g$  for various  $R_g$  at  $p_{T,jet}$  GeV/c



- Increasing  $p_{T,jet}$  has a mild effect on substructure
- Selection on  $R_g$  determines the  $z_g$  shape - high degree of correlation
- **Phase space restrictions have a large impact!**

$$\#\# < p_{T,\text{jet}} < \#\# \text{ GeV}/c$$

$$\#\# < p_{T,\text{initiator}} < \#\# \text{ GeV}/c$$

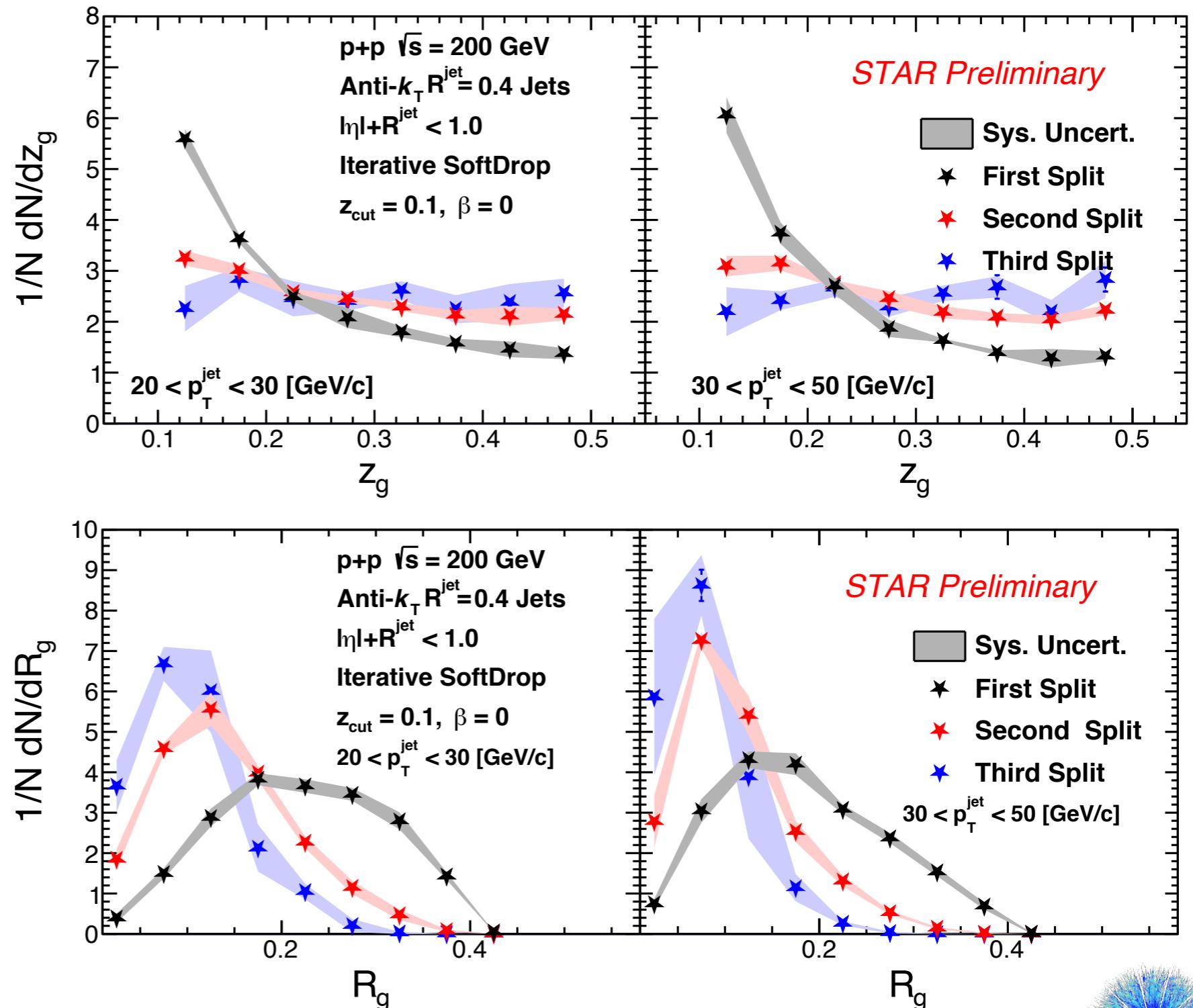


Evolution of the splitting  
observables as we travel  
along the jet shower

# Fully corrected results

1st, 2nd, 3rd splits for various  $p_{T,jet}$

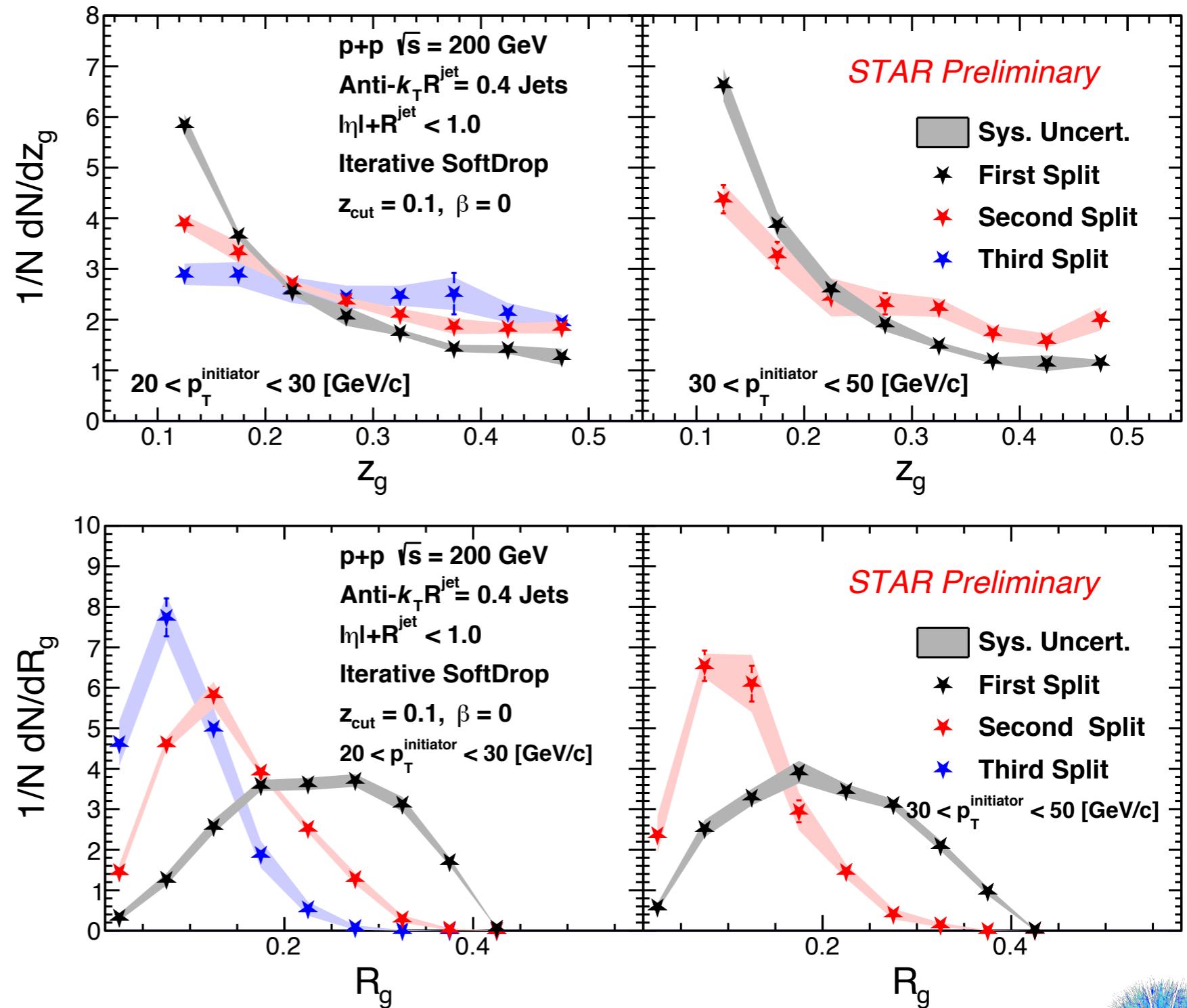
- For a given  $p_{T,jet}$ , what are the  $z_g, R_g$  at 1st, 2nd and 3rd splits? Follow a jet...
- Significant differences between **first**, **second** and **third** splits
- **Flat  $z_g$  distribution and smaller  $\langle R_g \rangle$**  for the third split, where we observe collinear emissions



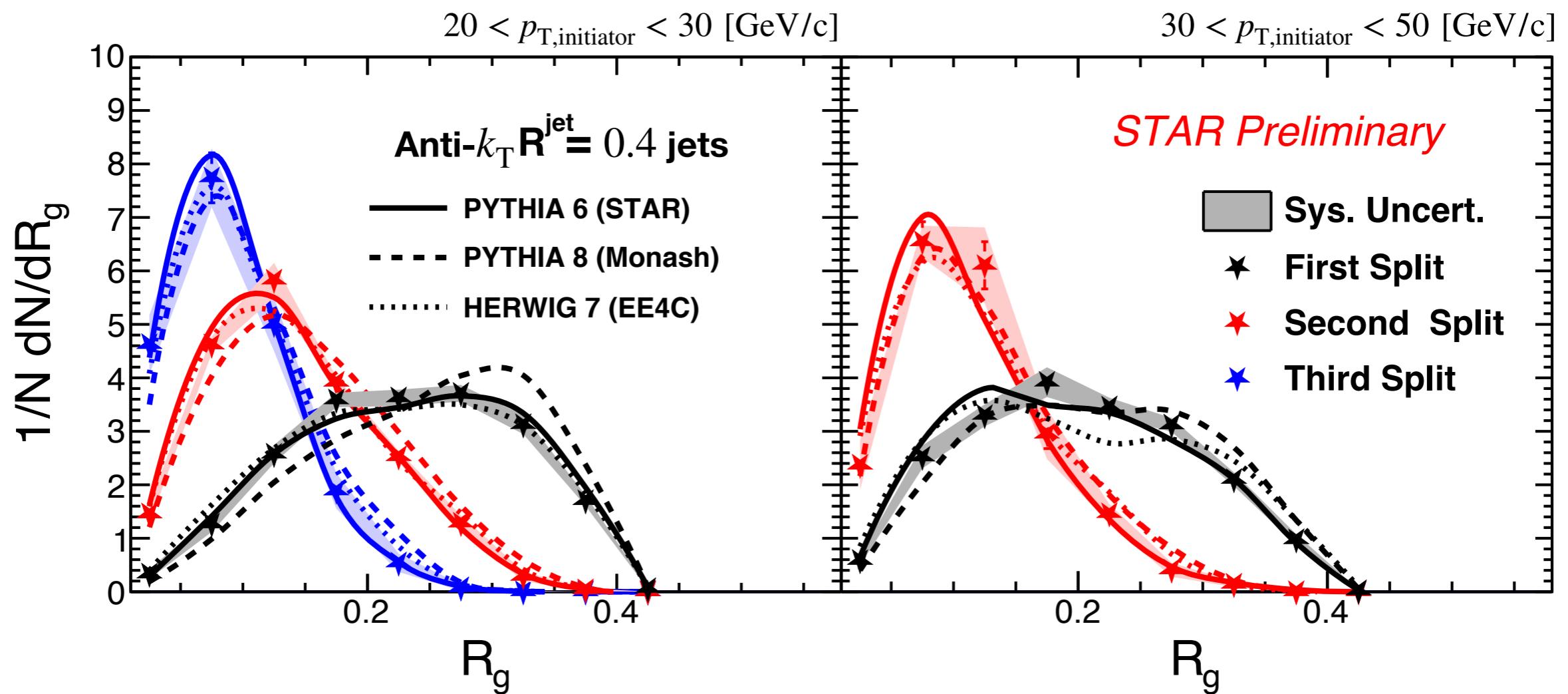
# Fully corrected results

## 1st, 2nd, 3rd splits for various $p_{T,\text{initiator}}$

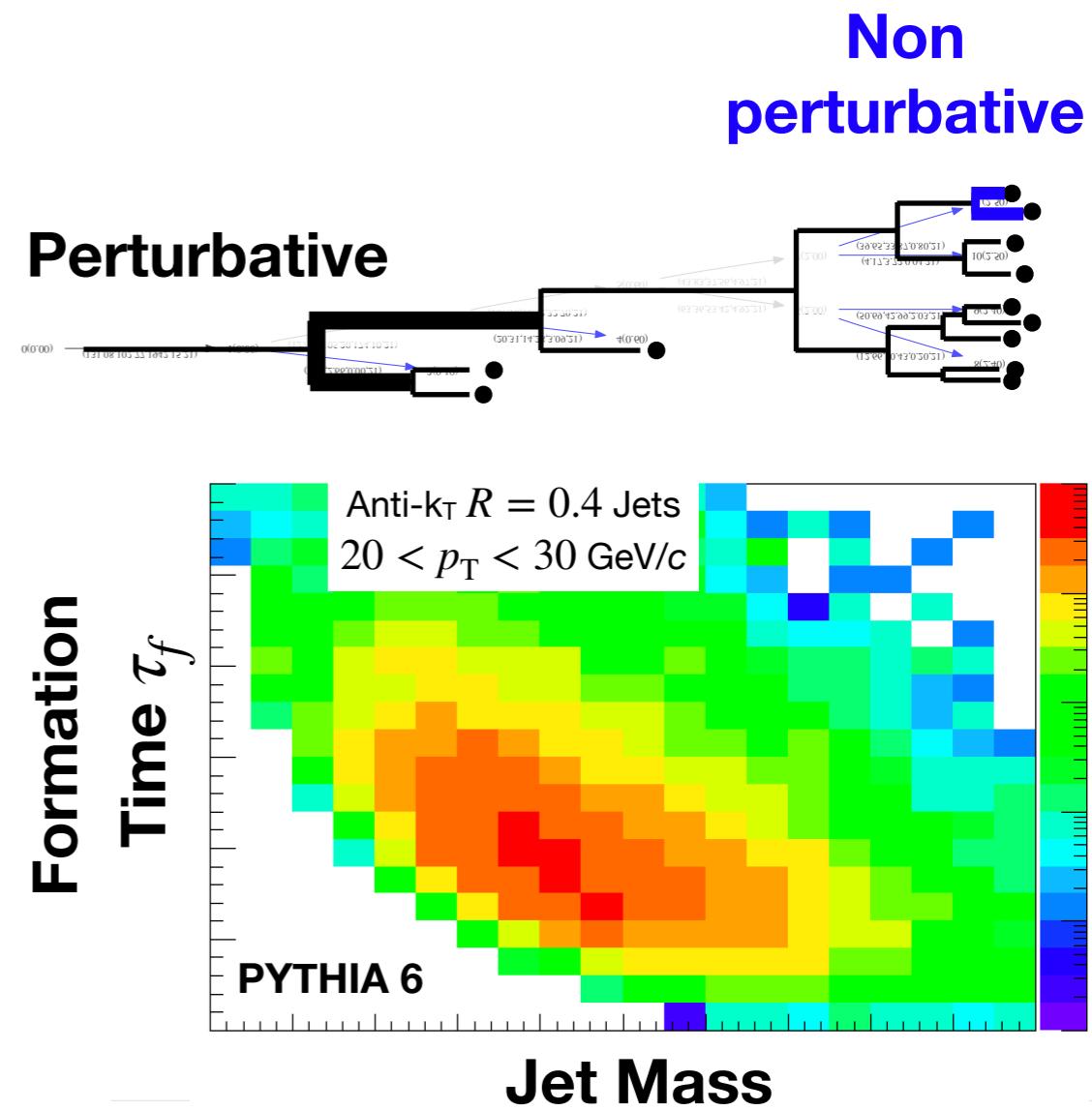
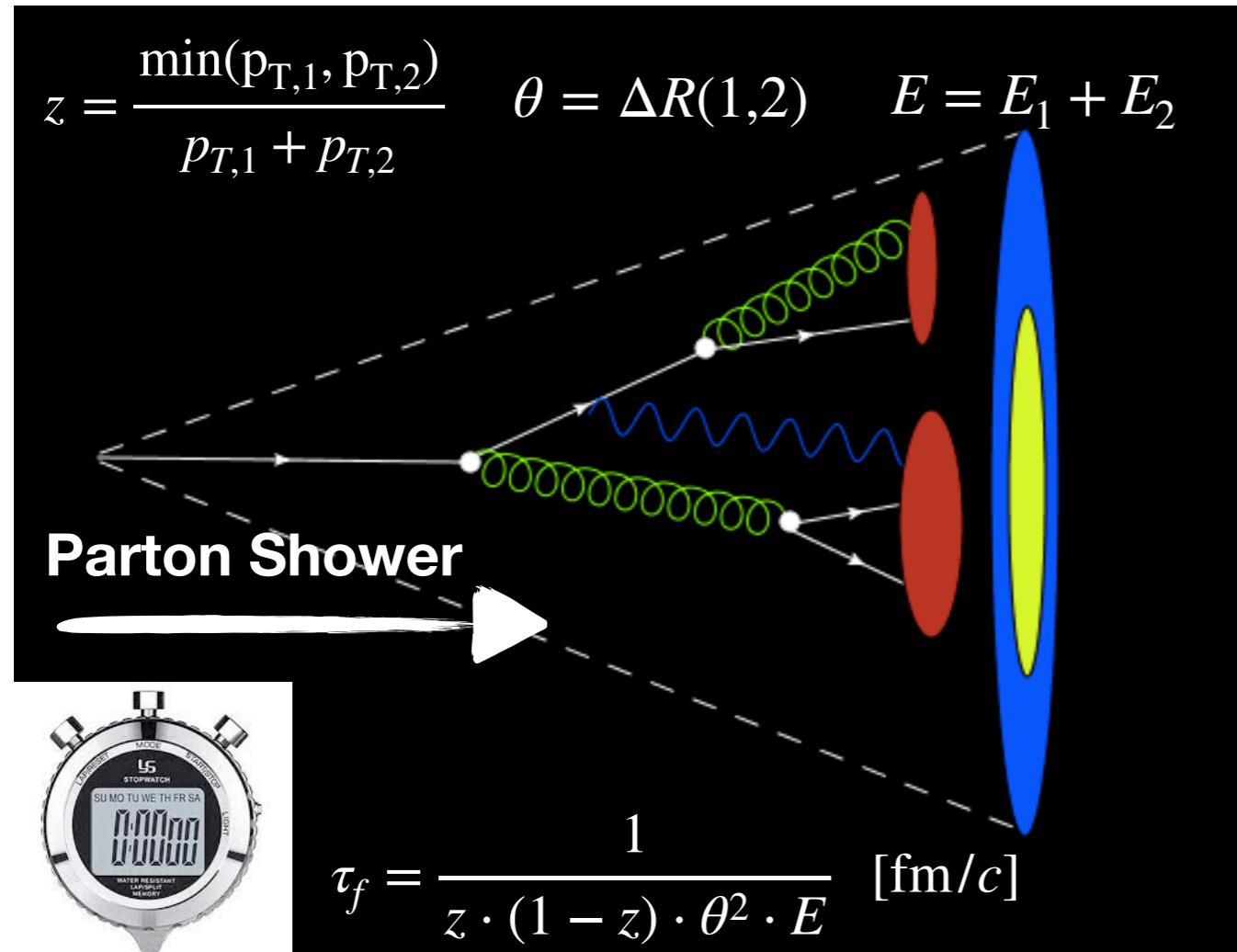
- For a given split with  $p_{T,\text{initiator}}$ , what are the  $z_g, R_g$  for 1st, 2nd and 3rd splits? Follow a split...
- Splits are directly comparable to each other - only difference is where they occur in the shower
- Hint of differences at the **second split  $z_g$**  (similar  $R_g$ ) for initiator vs. jet momenta selection



# Comparisons with leading order MC - $R_g$ for various $p_{T,\text{initiator}}$



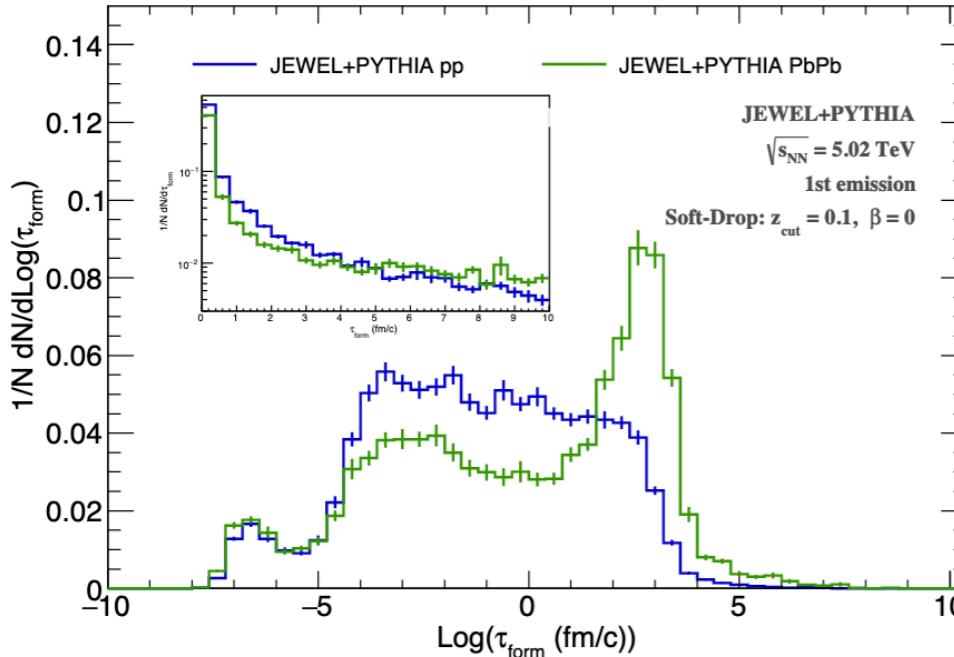
- Three MC (PYTHIA 6, PYTHIA 8, HERWIG 7) models describe the overall trend of narrowing of jet substructure for higher splits
- Availability of emission phase space depends on both jet momenta and split number - similar peaks of  $R_g$  for **third splits** on the left to **second splits** on the right



Formation time connecting  
jet splitting and charged  
particles in the jet

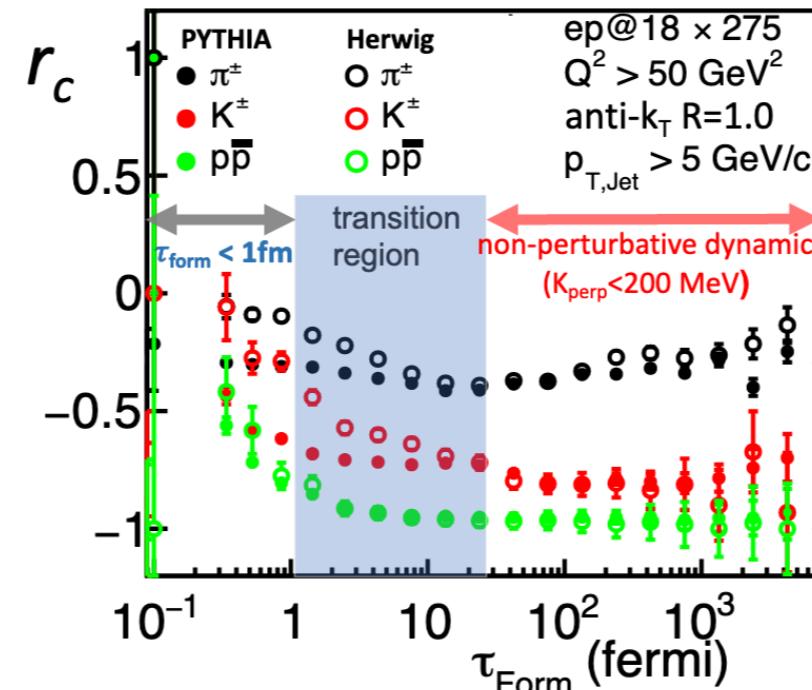
# Recent work on formation times

Apolinario et. al. *Eur.Phys.J.C* 81 (2021) 6, 561



- $\tau_f$  is a combination of substructure observables that results in a ‘time’
- Ensemble distributions of  $\tau$  contain information related to the parton shower
- Useful handle in jet quenching studies

- Recent studies also show its usefulness from the theoretical POV on isolating regions where calculations are valid
- Fuzzy area, but overall one can separate out ‘**mostly perturbative** and ‘**mostly non-perturbative** regions based on  $\tau$



$$r_c \equiv \frac{N_{CC} - N_{C\bar{C}}}{N_{CC} + N_{C\bar{C}}}$$

**Formation time :**  $[2z(1-z) P]/k_{\text{perp}}^2$   
 $z$  : momentum fraction of NL particle  
 $k_{\text{perp}}$ : relative transverse momentum between L & NL

- There is strong flavor dependence in  $r_c$
- In specific kinematic region PYTHIA and Herwig differ significantly

Mondal et. al. @ BOOST 2021

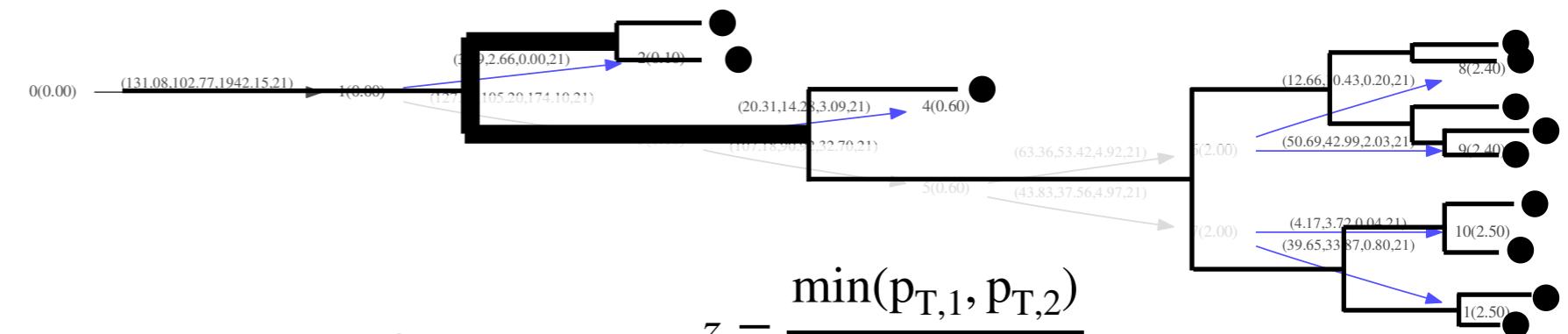
# Formation time in jets

Apolinario et. al. *Eur.Phys.J.C* 81 (2021) 6, 561

- SoftDrop  
first split  $\tau_f$   
(varying  $z_{\text{cut}}$ )

$$\tau_f = \frac{1}{z \cdot (1 - z) \cdot \theta^2 \cdot E} \quad [\text{fm}/c]$$

Measure  $\ln \tau_f \in [-2,4]$

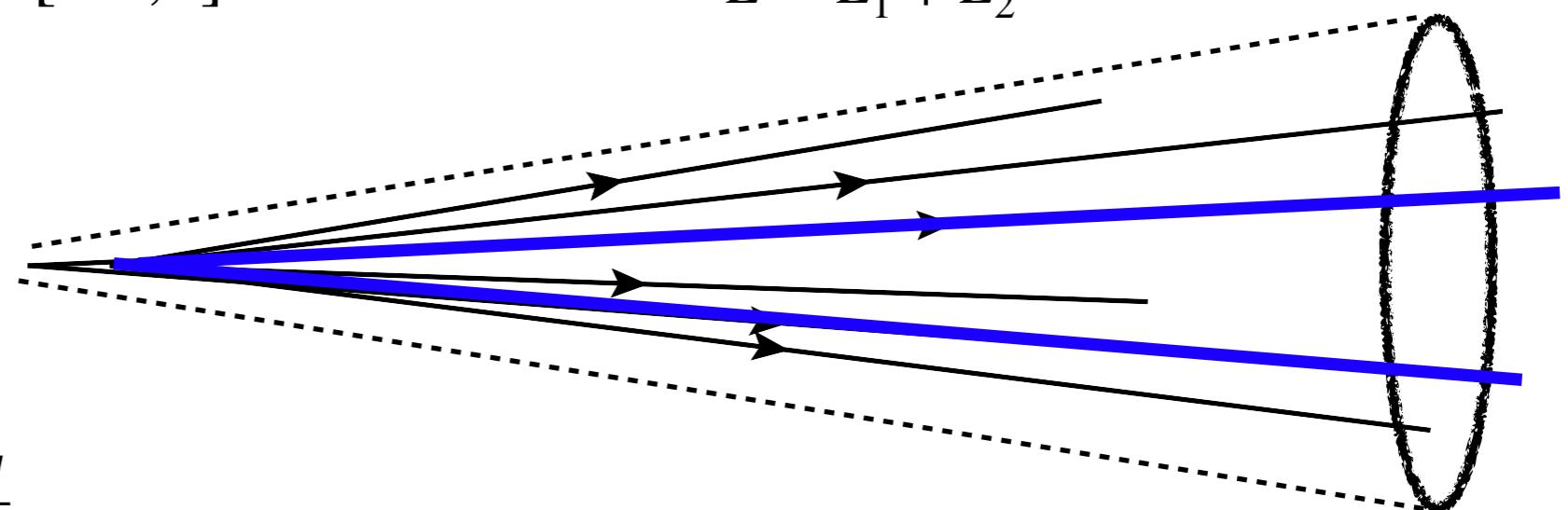


$$z = \frac{\min(p_{T,1}, p_{T,2})}{p_{T,1} + p_{T,2}}$$

$$\theta = \Delta R(1,2)$$

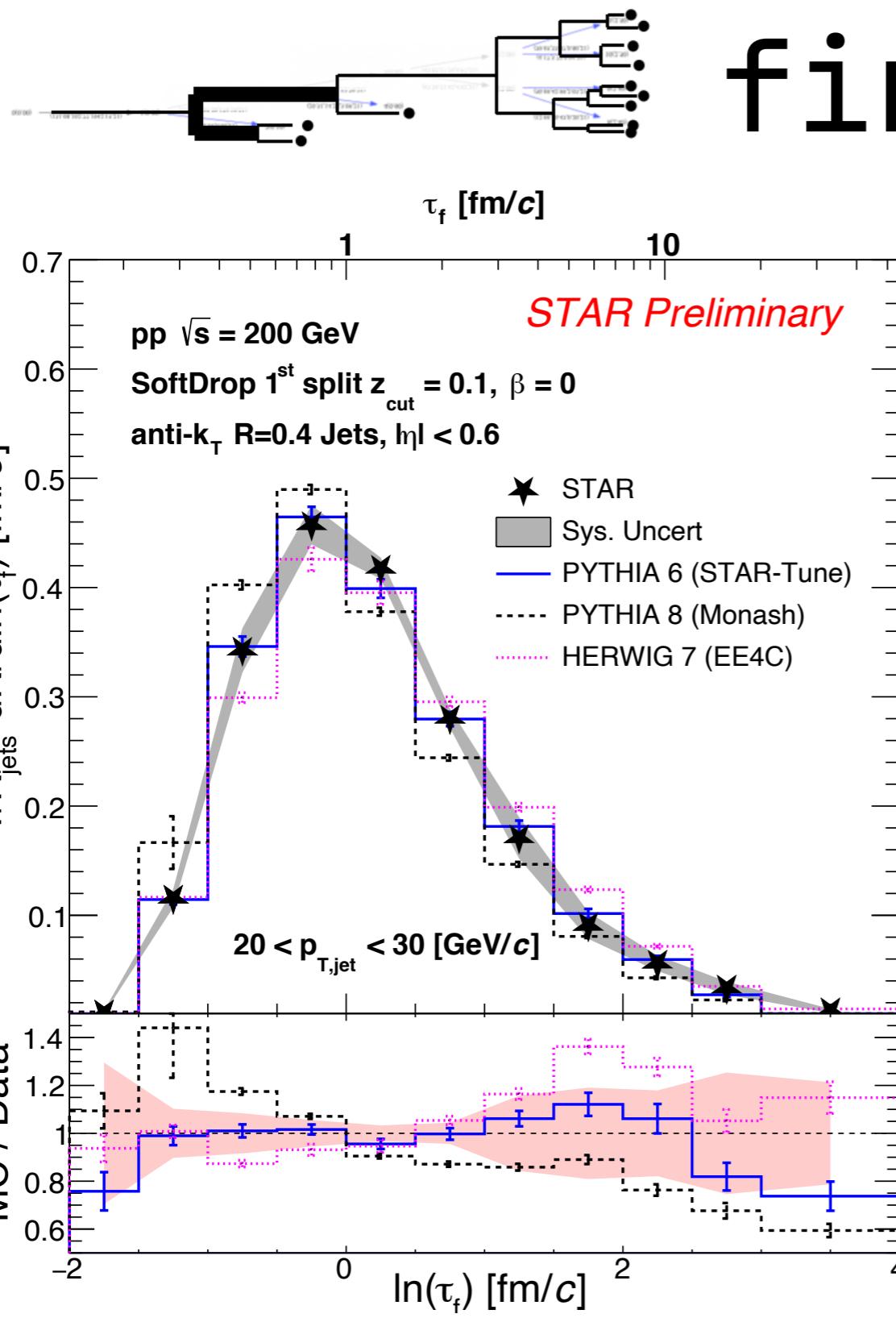
$$E = E_1 + E_2$$

- Leading and subleading ch-particle  $\tau_f$



Mondal et. al. @ BOOST 2021

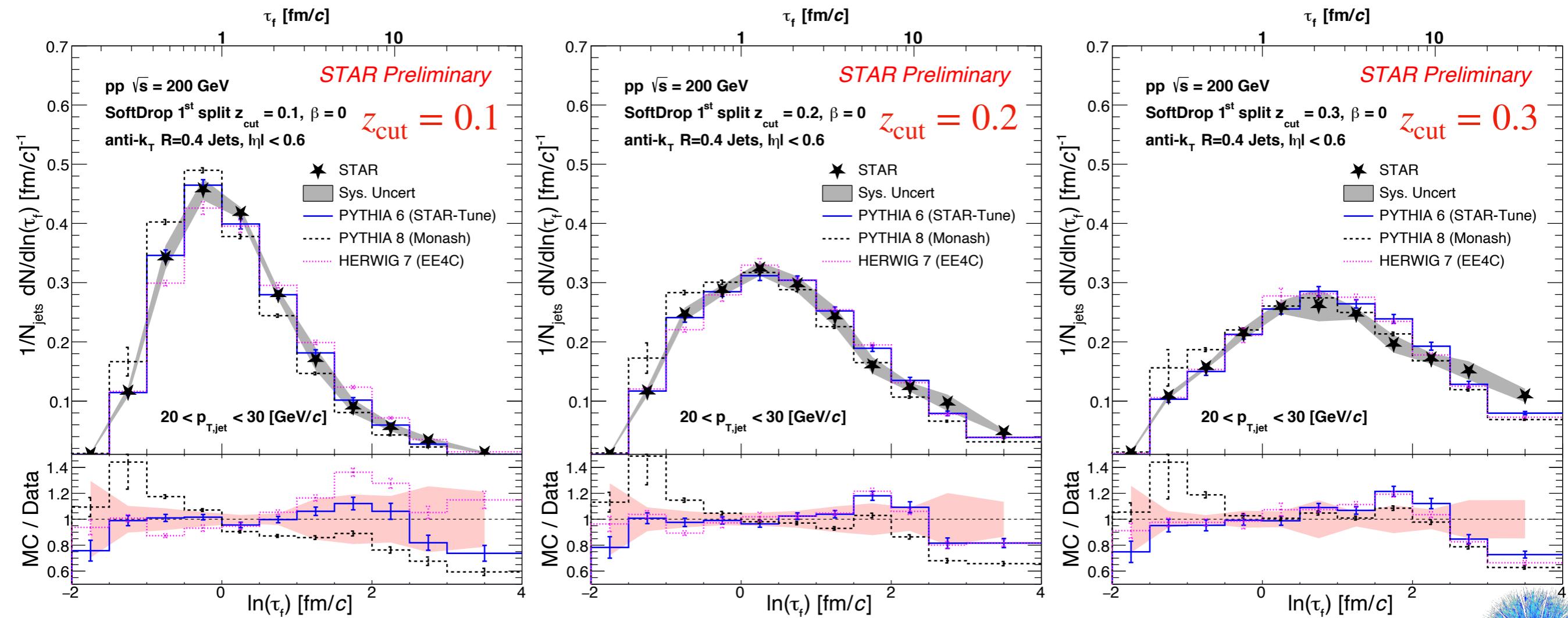
# Fully unfolded SoftDrop



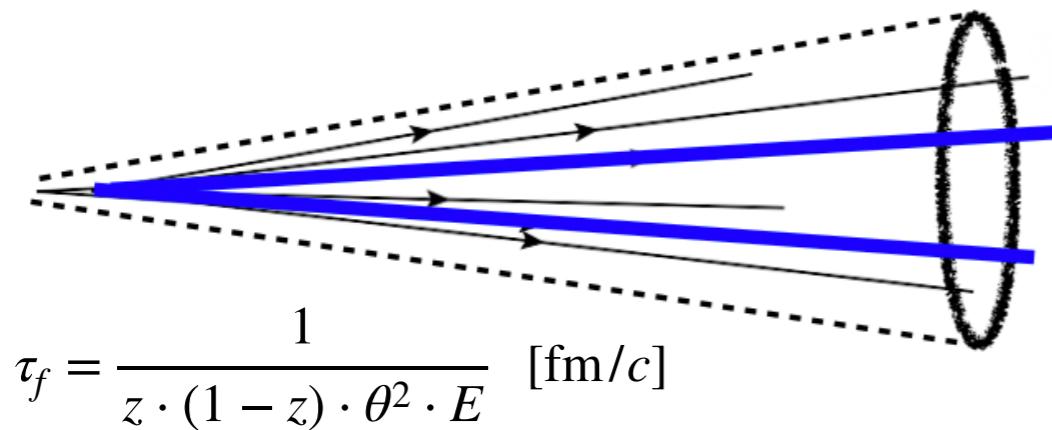
- Peak of the distribution sits roughly around 1 fm/c for jets in our kinematics with  $20 < p_T < 30$  [GeV/c]
- PYTHIA 6 generally describes the trend while HERWIG 7 and PYTHIA 8 sandwich the data
- How does this look like if we vary the SD criterion?

# Impact of increasing $z_{\text{cut}}$

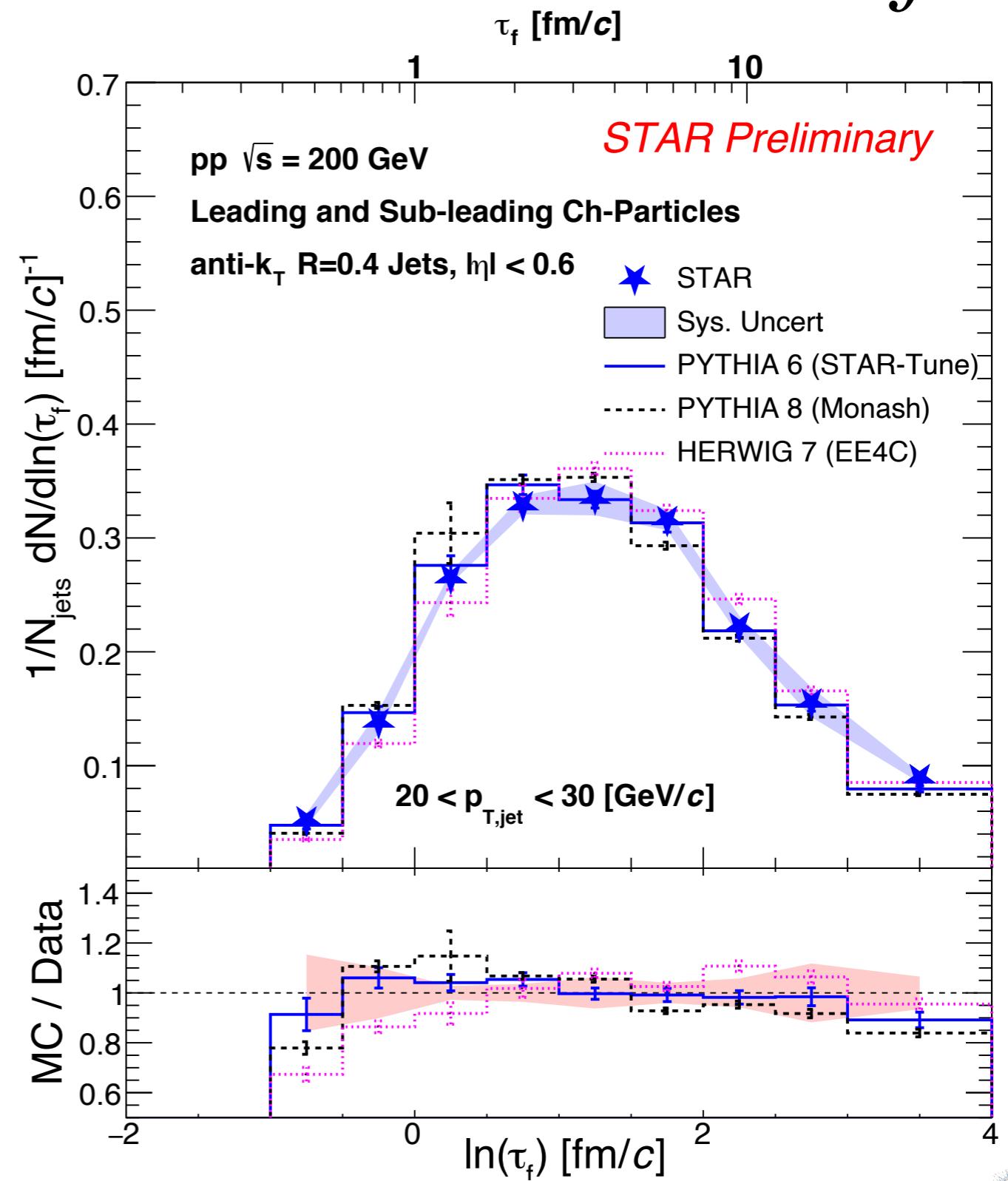
- Increasing  $z_{\text{cut}}$  effectively grooms away more branches before the criterion is passed
- Similar to traveling further along the shower - evident by the shift in the  $\langle \tau_f \rangle$  towards the right



# Charged particle $\tau_f$

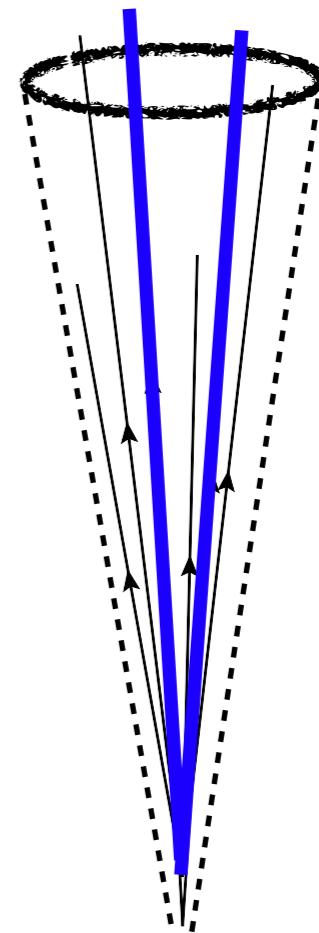
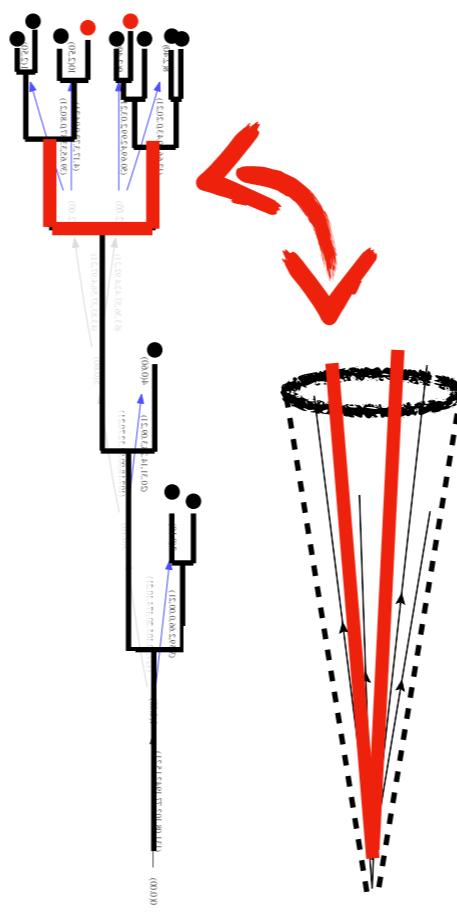
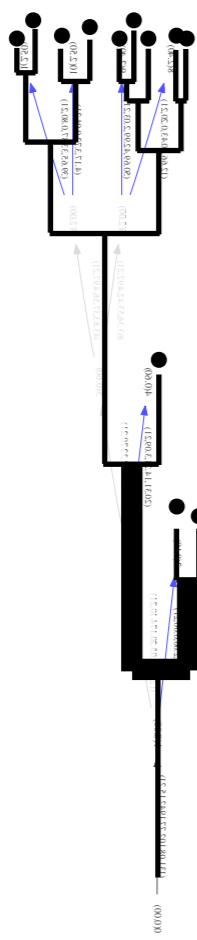


- $\tau_f$  calculated using the leading and sub-leading charged particles within jets
- These have much larger  $\langle \tau_f \rangle$  than the SD first splits
- These exist in the jet clustering tree somewhere - where/when exactly does the tree resolve these particles?

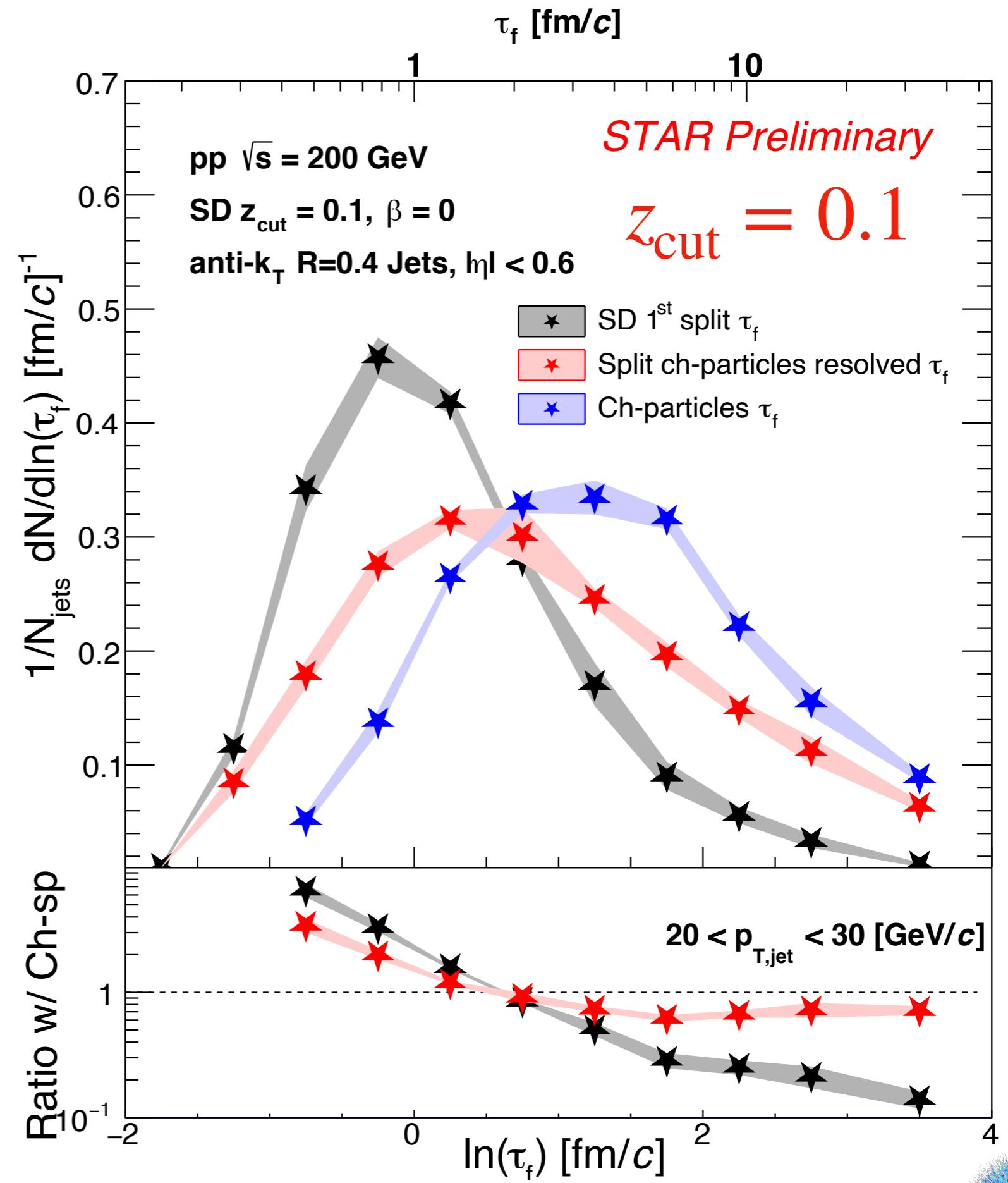


# Connecting the two regimes

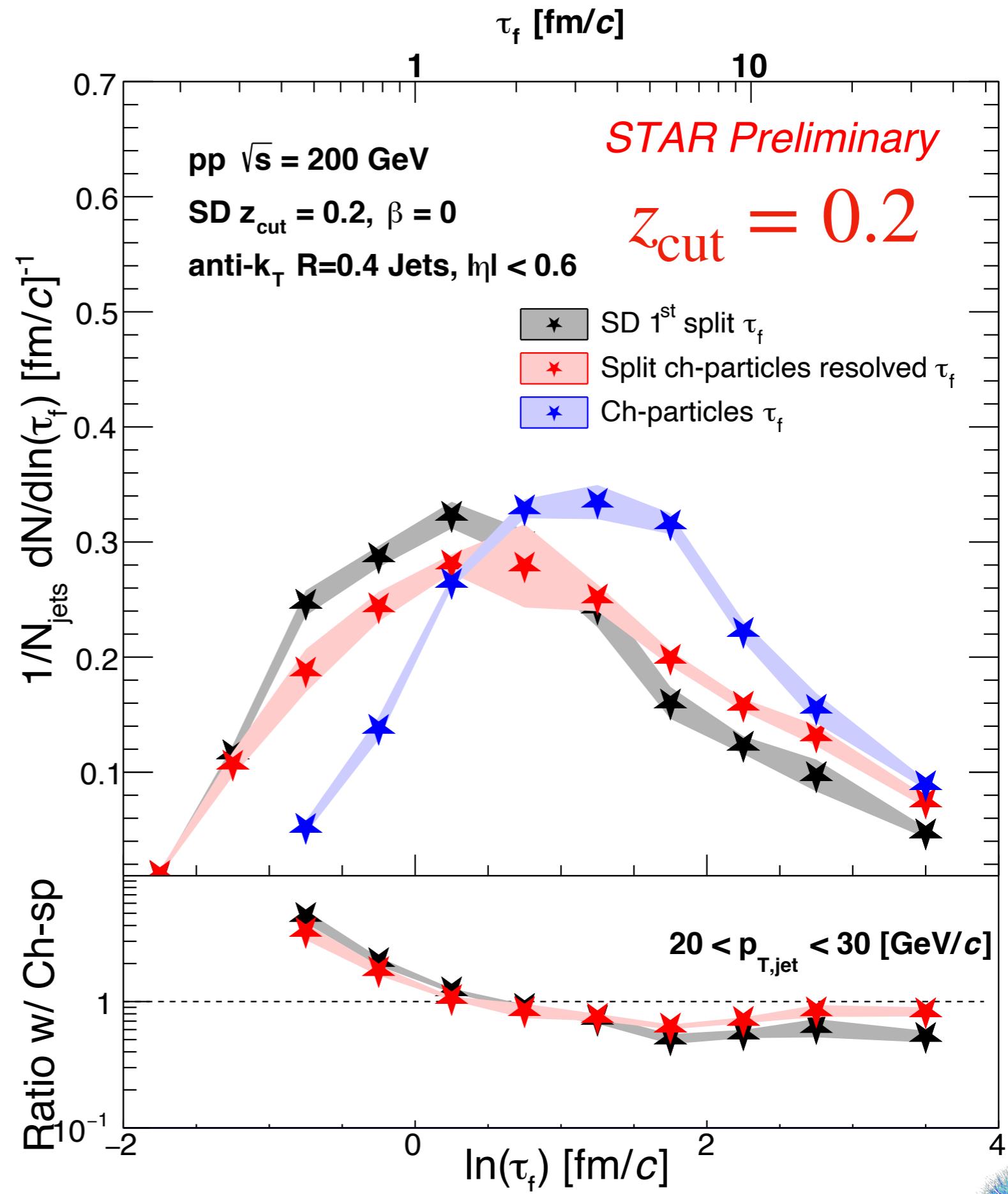
- SoftDrop first split  $\tau_f$  (varying  $z_{\text{cut}}$ )
- SoftDrop split (varying  $z_{\text{cut}}$ ) resolving the two leading charged particles
- Leading and subleading ch-particle  $\tau_f$



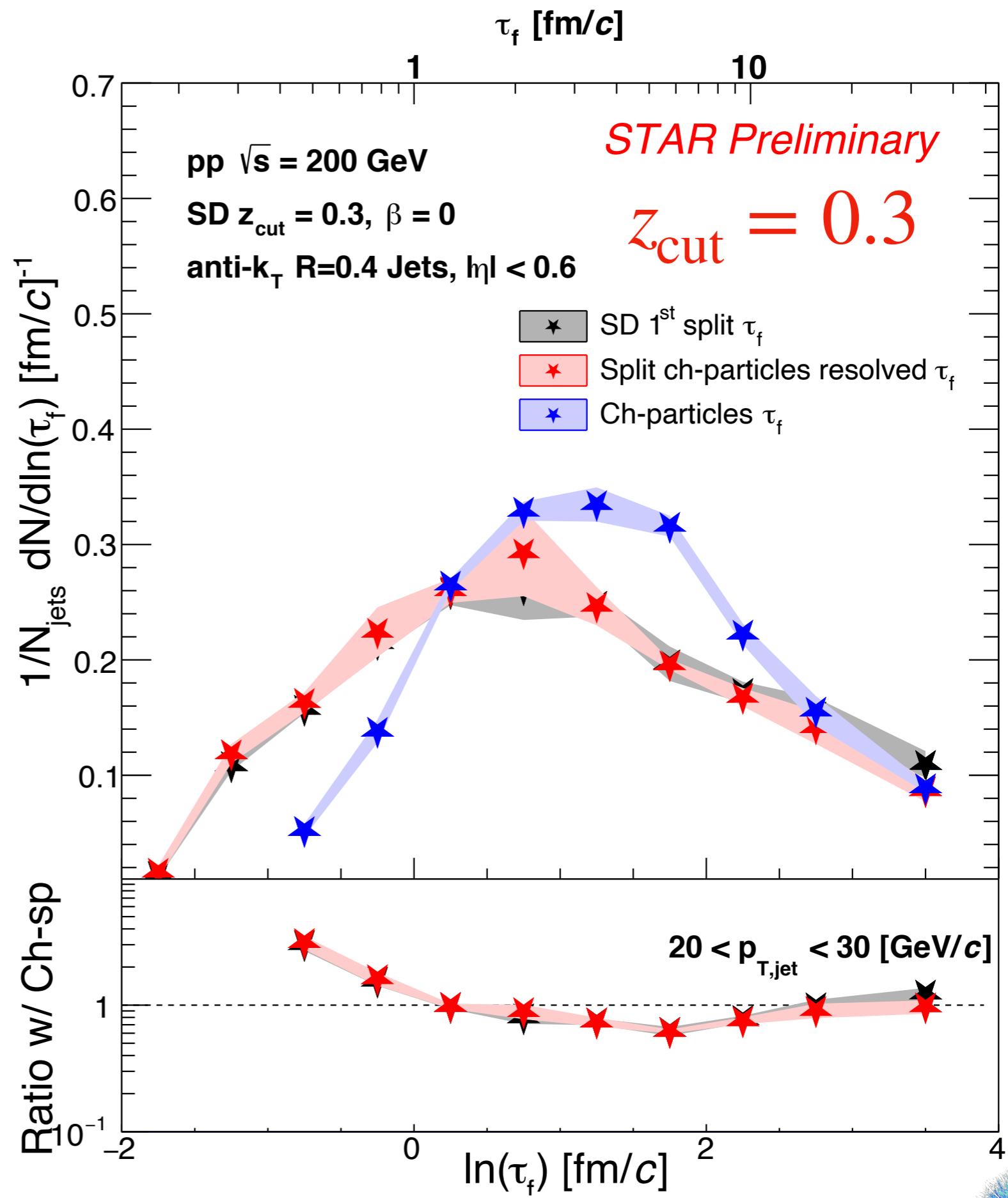
- SD first splits are significantly earlier in time compared to the charged particles
- At  $\approx 2 - 3$  fm/c, we see the ratio for SD splits go less than 1
- Resolved splits have similar shape to the ch-particle splits after  $\approx 3 - 5$  fm/c
- What happens if we increase  $z_{\text{cut}}$ ?



- Increasing to  $z_{\text{cut}} = 0.2$  requires harder splitting scale
- The **first split** and **resolved split** distributions start to become similar
- Meaning that we are moving **further along the shower in time** towards the **charged particle  $\tau_f$**

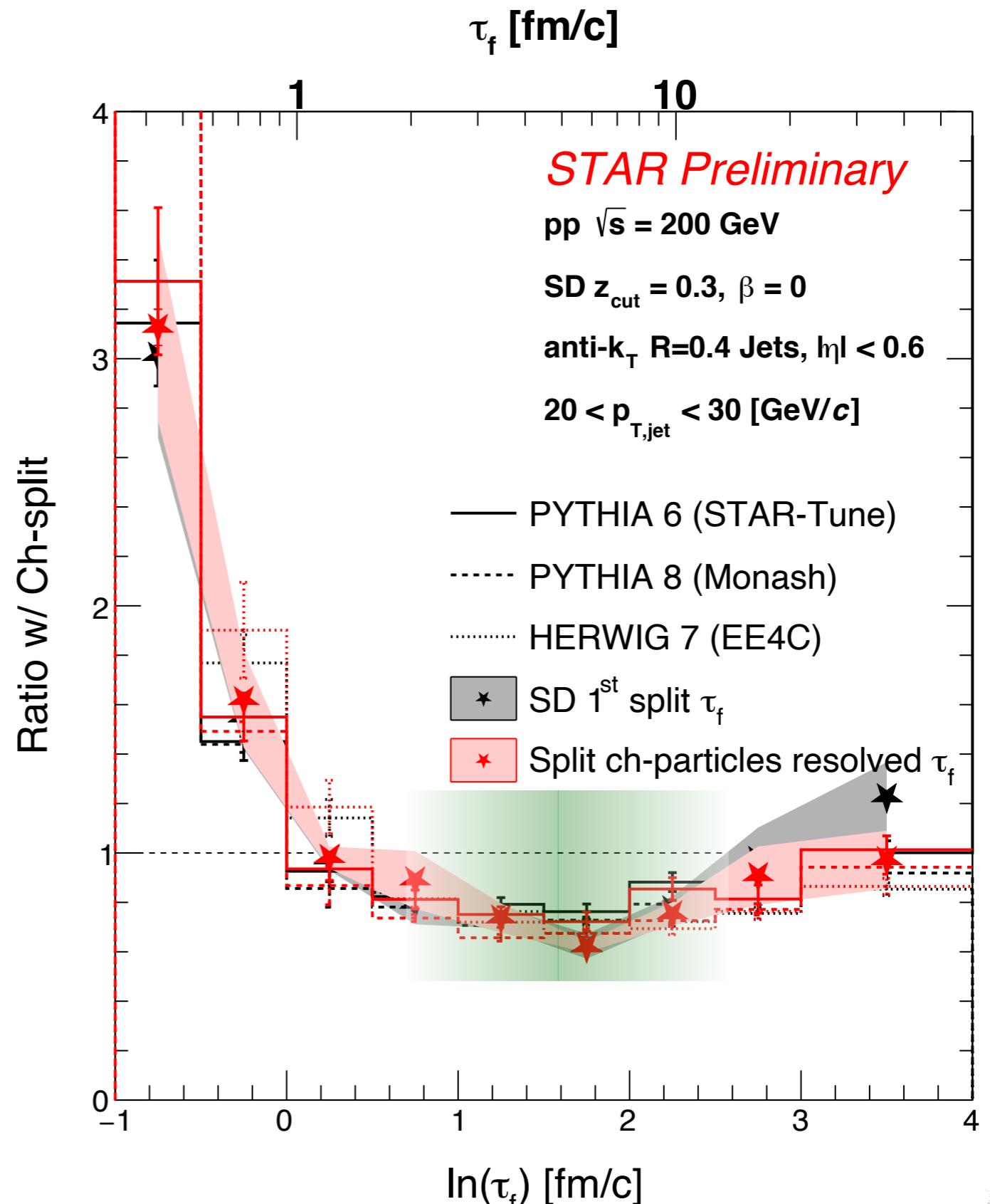


- No difference between the first spilt with extreme  $z_{\text{cut}}$  and the resolved split!
- Surviving jets have a **formation time bias** to have later splits that are in the **(mostly) non-perturbative region**



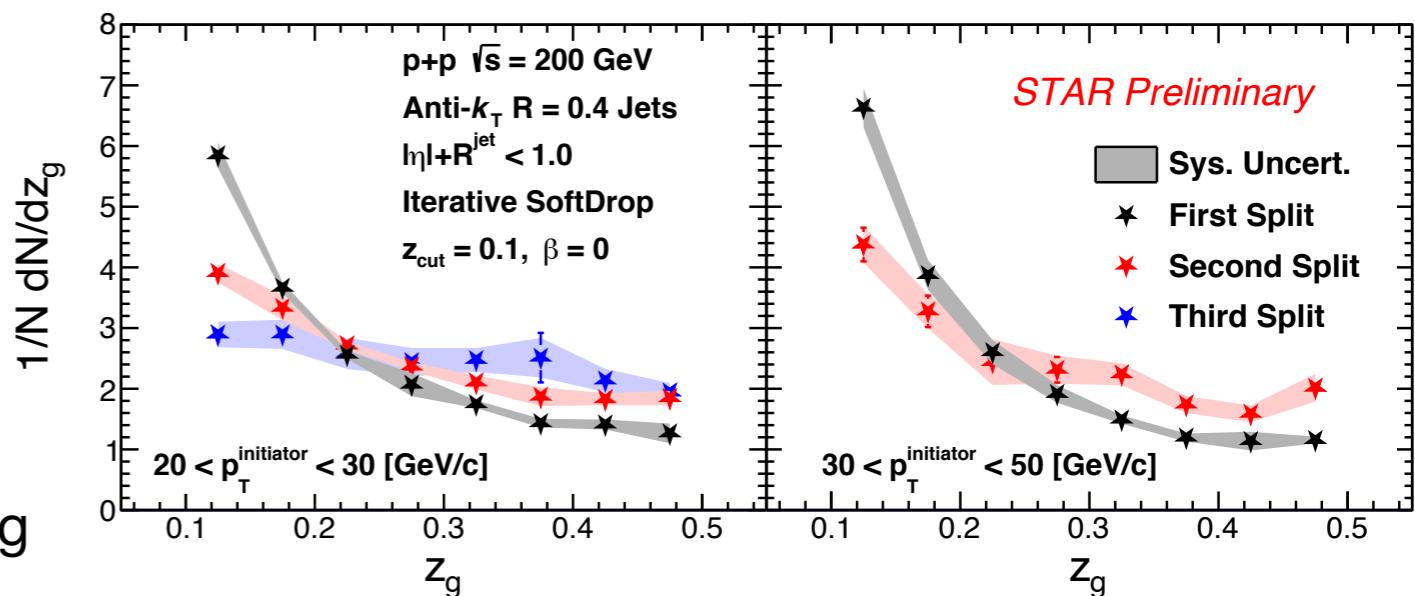
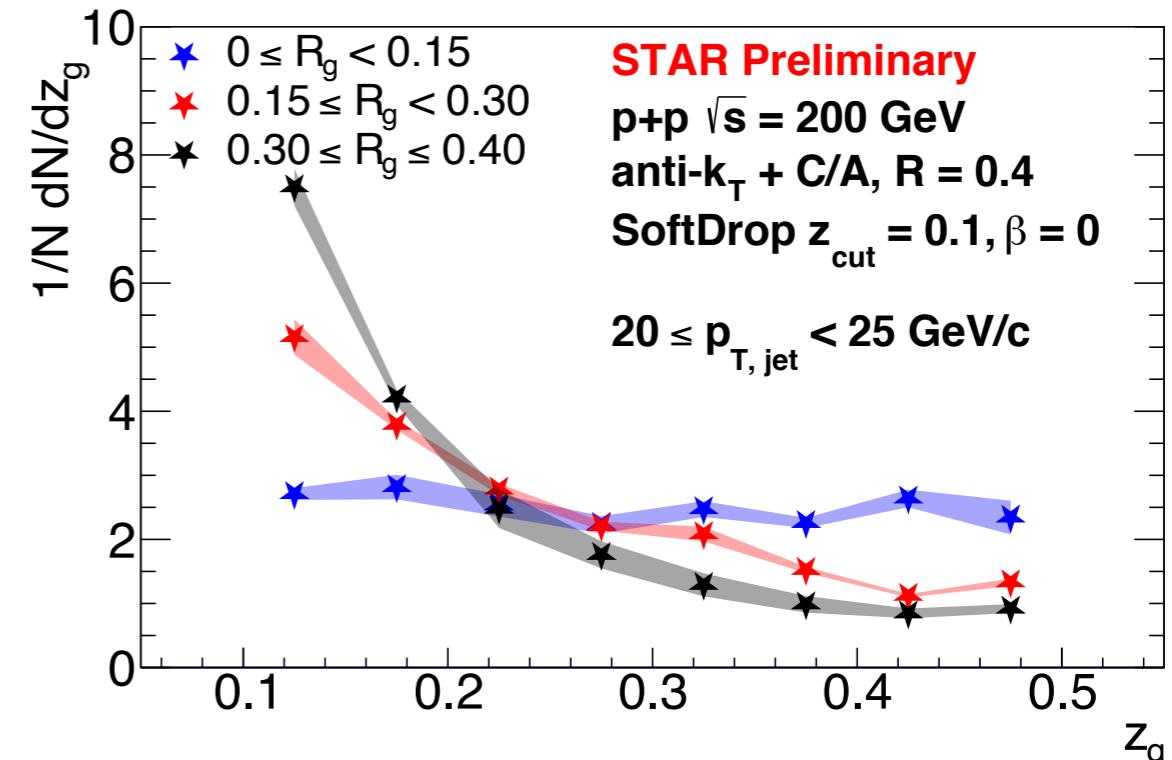
# Comparison to MC

- PYTHIA 8, PYTHIA 6 and HERWIG 7 show similar behavior of crossover and flattening
- Hints of differences between PYTHIA 6/8 and HERWIG 7 in the crossover region (ratio goes from  $> 1$  to  $< 1$ )
- Comparison of hadronization models to data across the measured  $\tau_f$  range will be crucial in studying the pQCD  $\rightarrow$  npQCD transition region (shown in the green box on the figure)



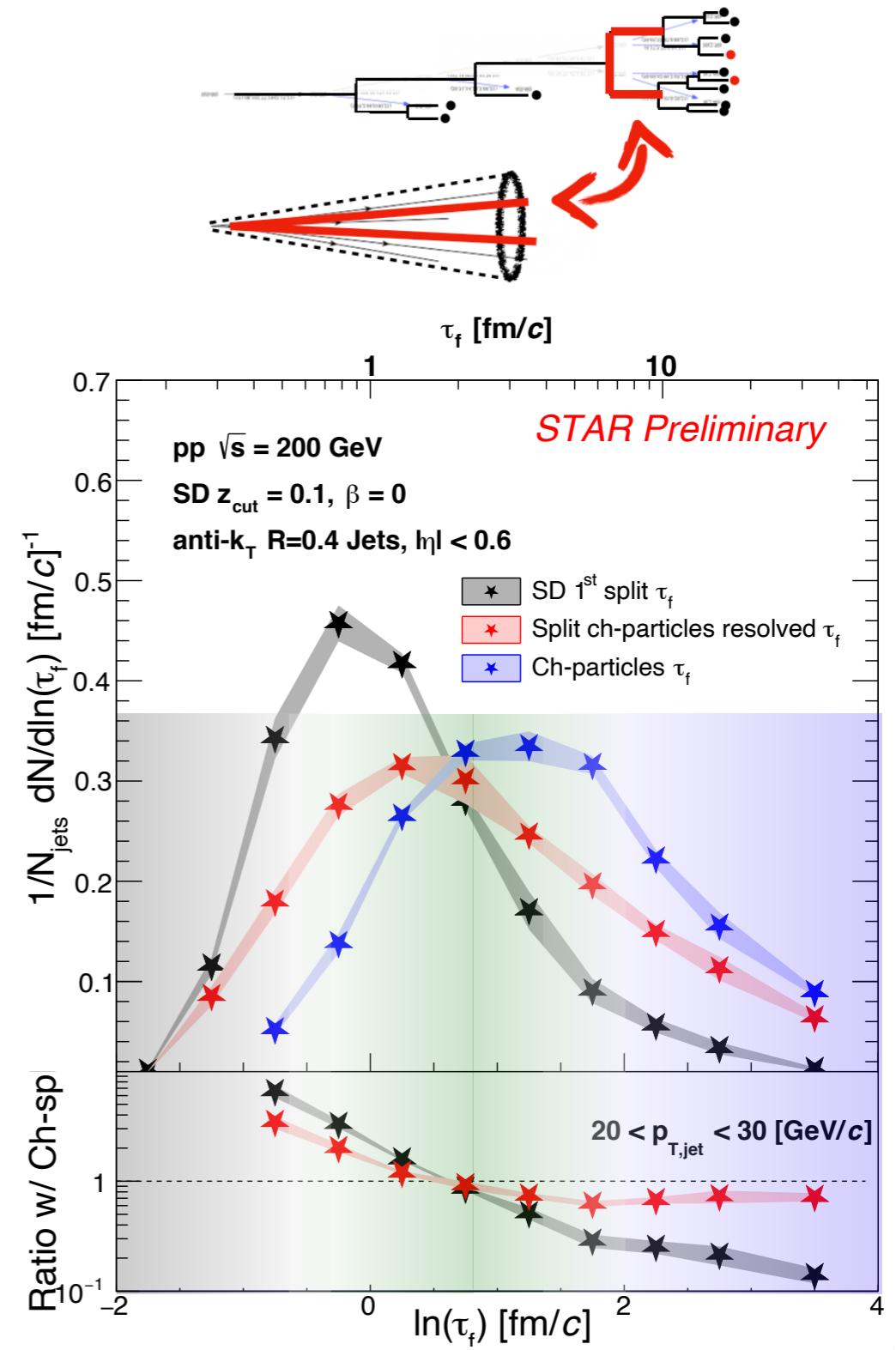
# Conclusions - I

- Jet substructure program at STAR aims at **mapping jet evolution** at RHIC energies
- Data show a **gradual variation in the available phase space**
  - leading to modifications (e.g. virtuality evolution) in the observed splitting kinematics
- Observe increased probability of **significantly harder/symmetric splittings** at the **third/narrow split** compared to the first and second splits
- Subjets at RHIC allow to **disentangle perturbative and non-perturbative dynamics of jet evolution** - these **third and narrow splits** for our low  $p_T$  jets end up being quite close to the  $\Lambda_{\text{QCD}}$  scale



# Conclusions - II

- First measurements of formation time from the jet splitting trees and from charged particles in the jet
- As the  $z_{\text{cut}}$  increases, the **crossover in  $\tau_f$**  moves to the earlier time
- Resolved SD splits show similar shape as the **charged particle split** at large  $\tau_f$  values occurring in the predominantly **non-perturbative region**
- Comparison of the different splits highlights the transition from **pQCD** to **npQCD**



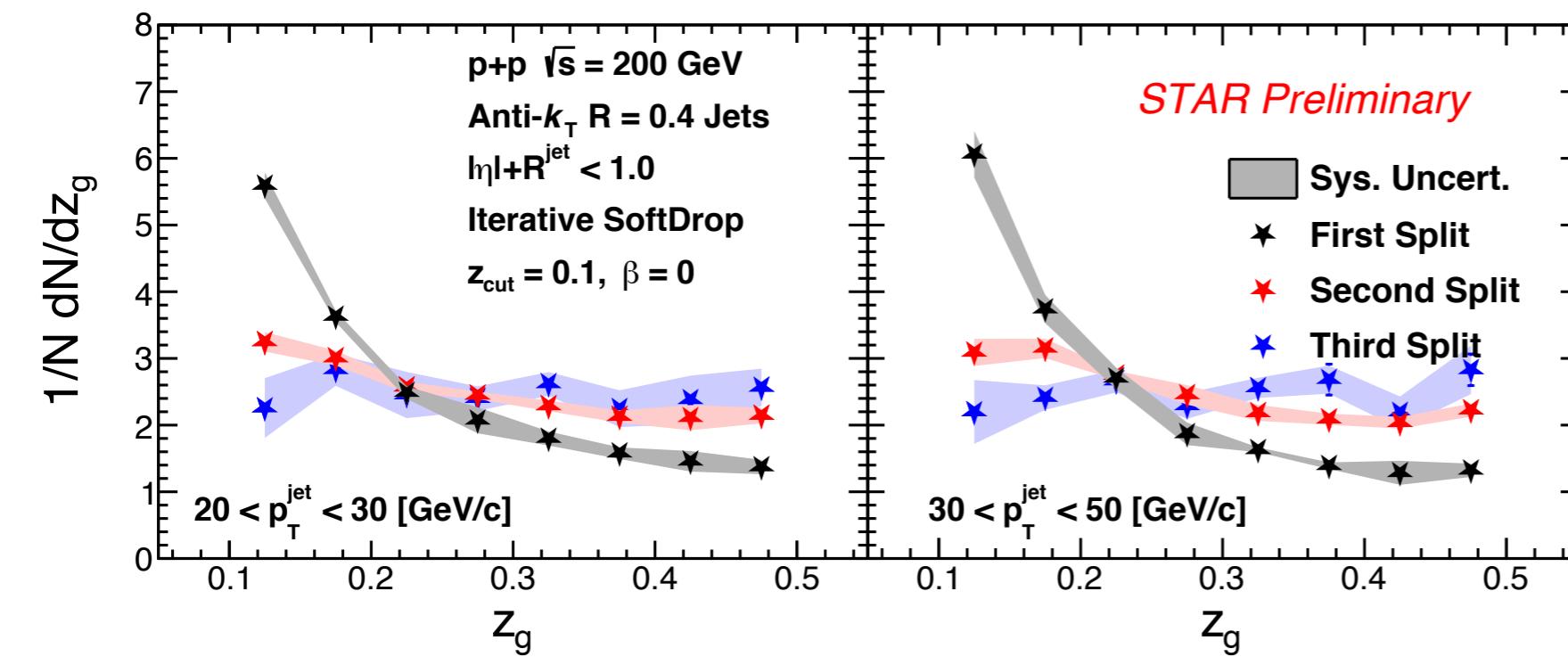
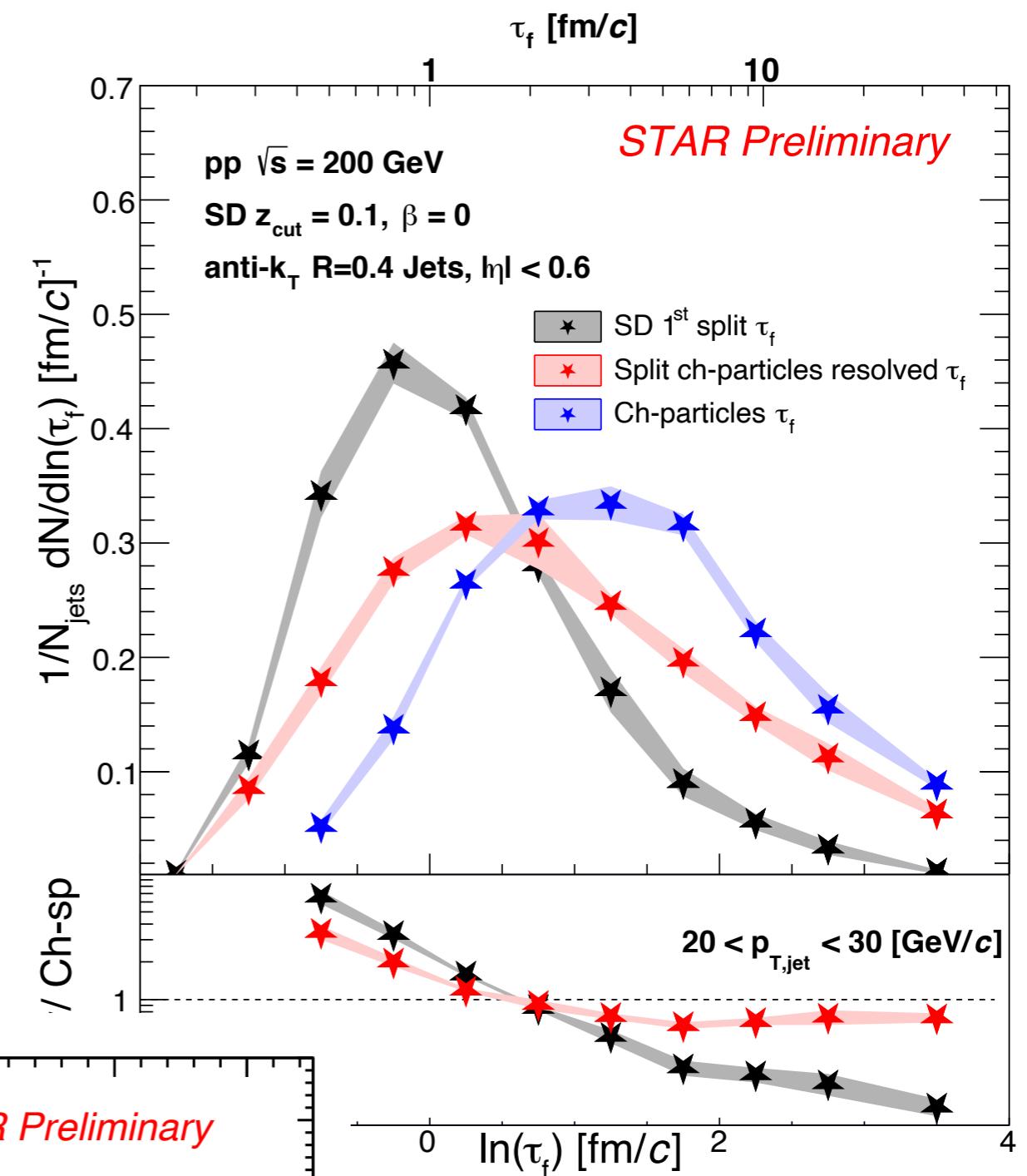
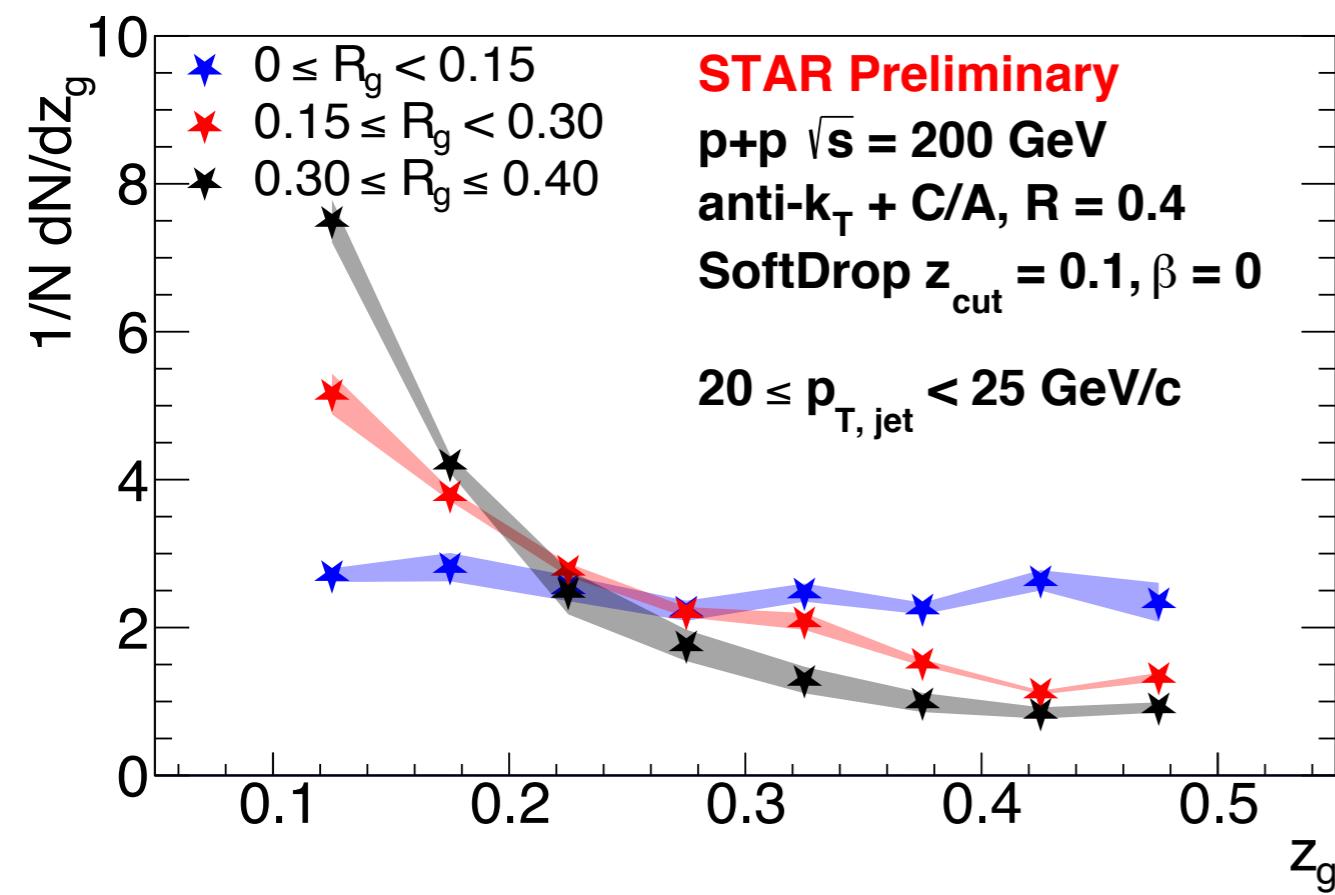
# Backup

# Next steps

- In our upcoming final results we will delve further into comparisons
  1. Various handles in the MC -

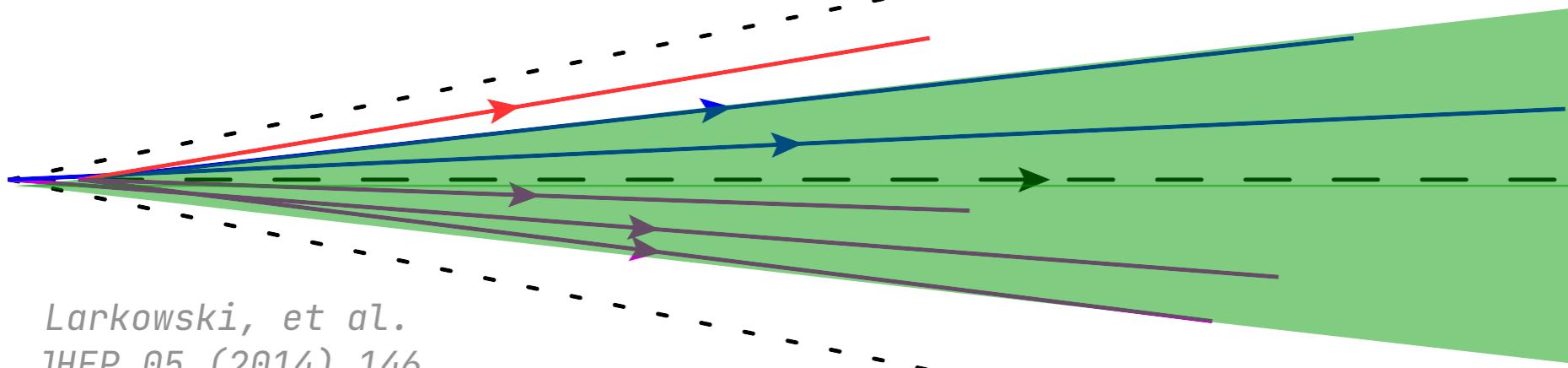
	Parton Shower	Hadronization	UE Tune
PYTHIA-6	$k_T$	Lund	RHIC
PYTHIA-8	Dipole/Vincia/ $p_T$	Lund	RHIC/LHC
HERWIG-7	Angular/Dipole	Cluster	LHC
SHERPA	Dipole	Lund/Cluster	LHC

2. In discussion with our theory colleagues on feasibility of calculations
- First measurement that can potentially **distinguish experimental quantities according to a ‘time scale’** via formation time of splits
    - Extremely useful in a heavy ion environment



# SoftDrop

## Grooming criterion

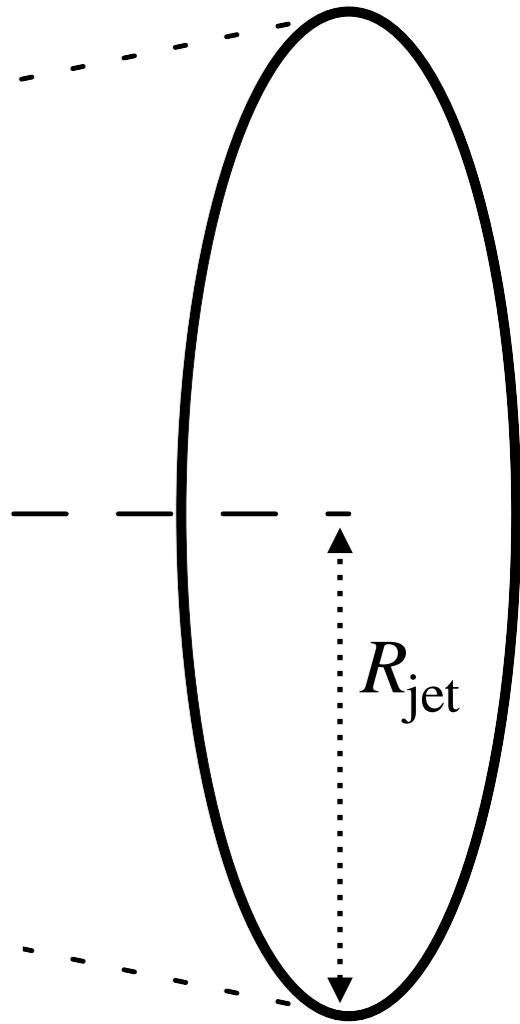
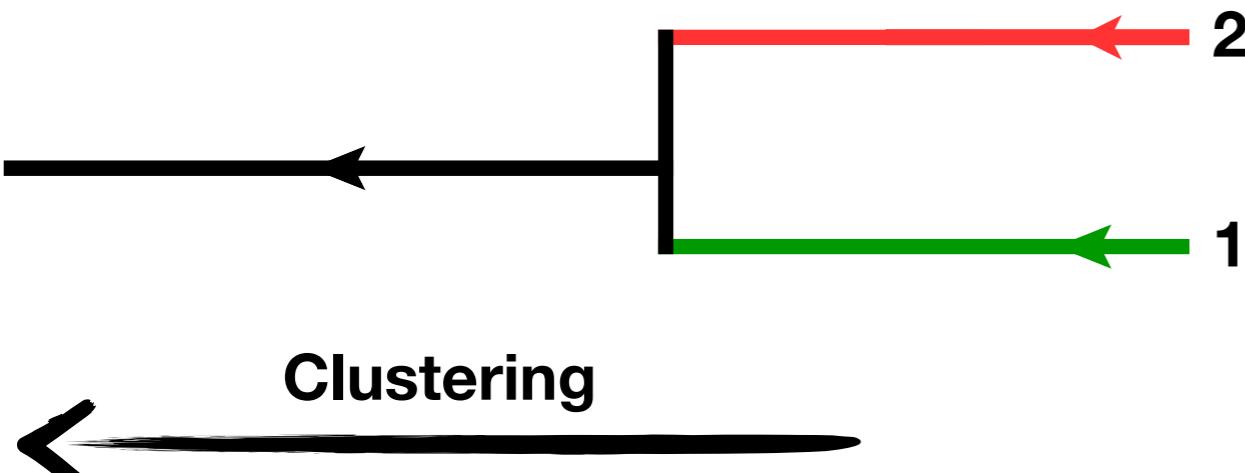


Larkowski, et al.  
JHEP 05 (2014) 146  
Dasgupta et al.  
JHEP 09 (2013) 029

- Require subjet momentum fraction to pass

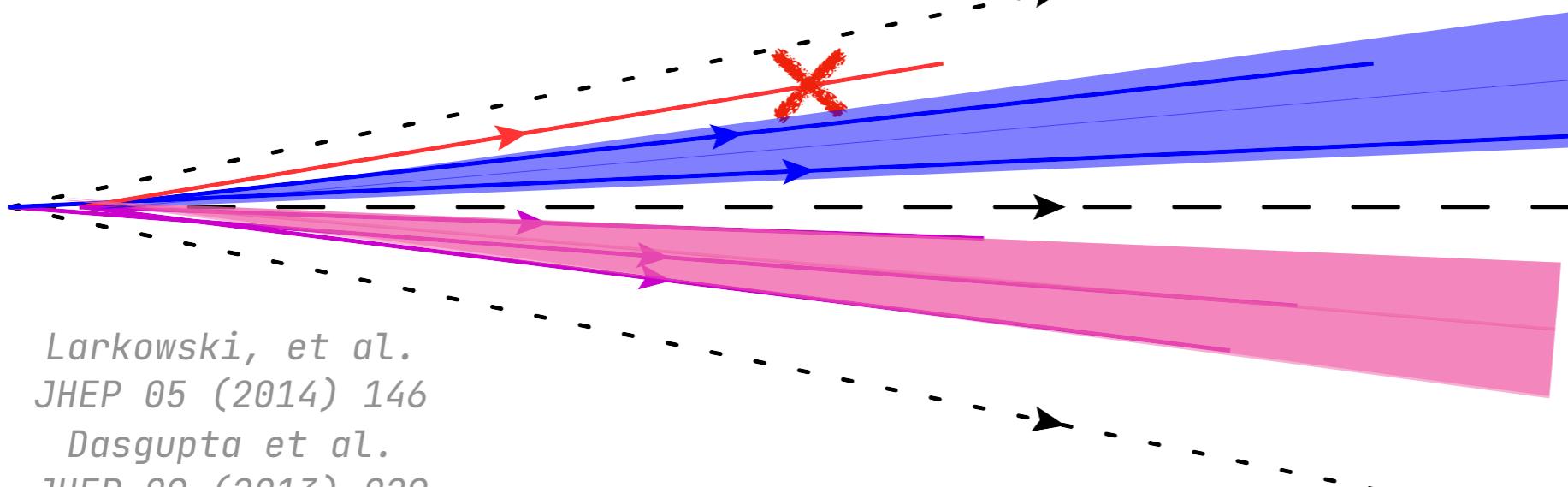
$$z_g = \frac{\min(p_{T,1}, p_{T,2})}{p_{T,1} + p_{T,2}} > z_{\text{cut}}(R_g/R_{\text{jet}})^{\beta}$$

$z_{\text{cut}} = 0.1$   
 $\beta = 0$



# SoftDrop

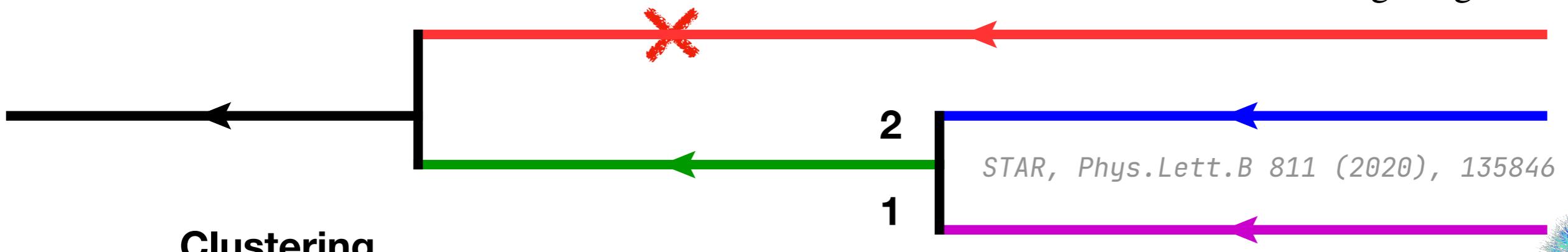
## Grooming criterion



- Require subjet momentum fraction to pass

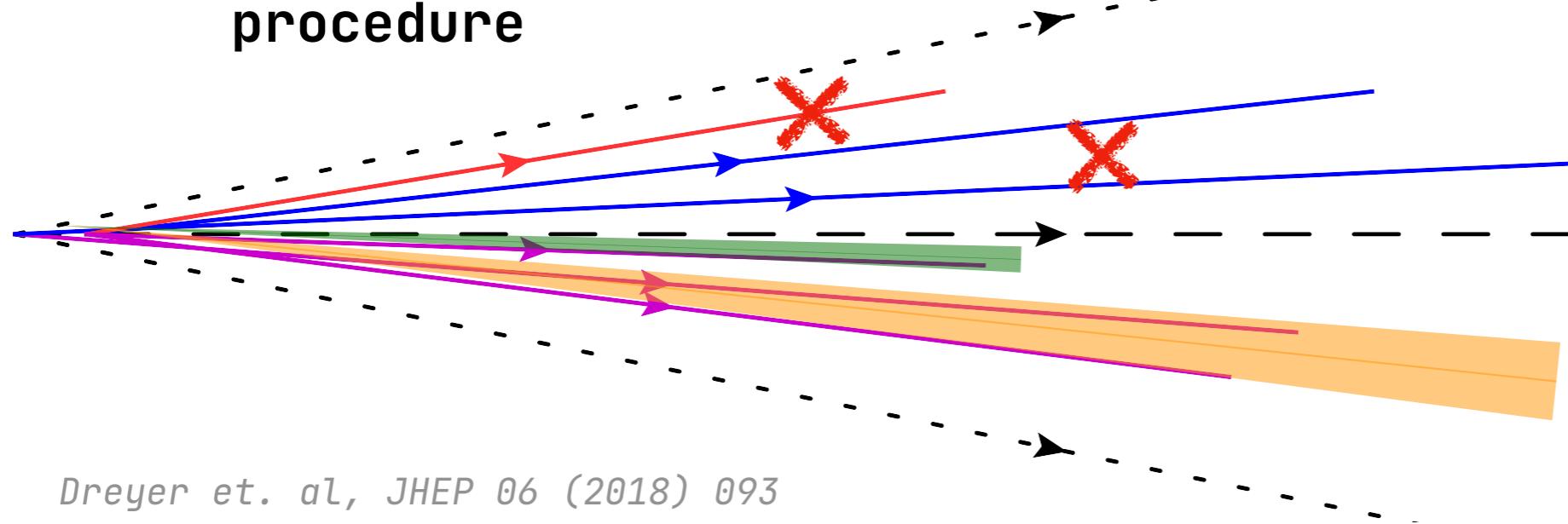
$$z_g = \frac{\min(p_{T,1}, p_{T,2})}{p_{T,1} + p_{T,2}} > z_{\text{cut}}(R_g/R_{\text{jet}})^{\beta}$$
$$z_{\text{cut}} = 0.1$$
$$\beta = 0$$

- With the two surviving branches (first hard split) - we define observables that characterize jet substructure  $z_g, R_g$



# SoftDrop

Extending the procedure

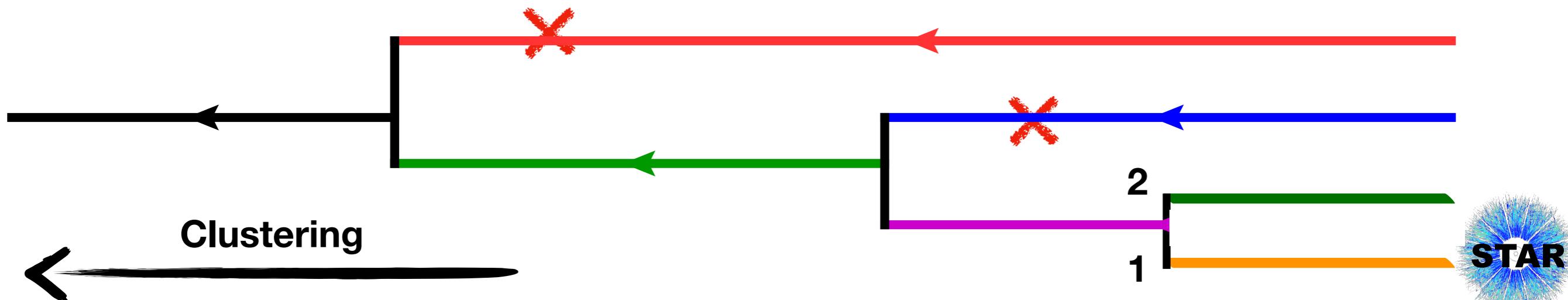


Dreyer et. al, JHEP 06 (2018) 093

We can implement the SoftDrop procedure throughout the C/A tree -

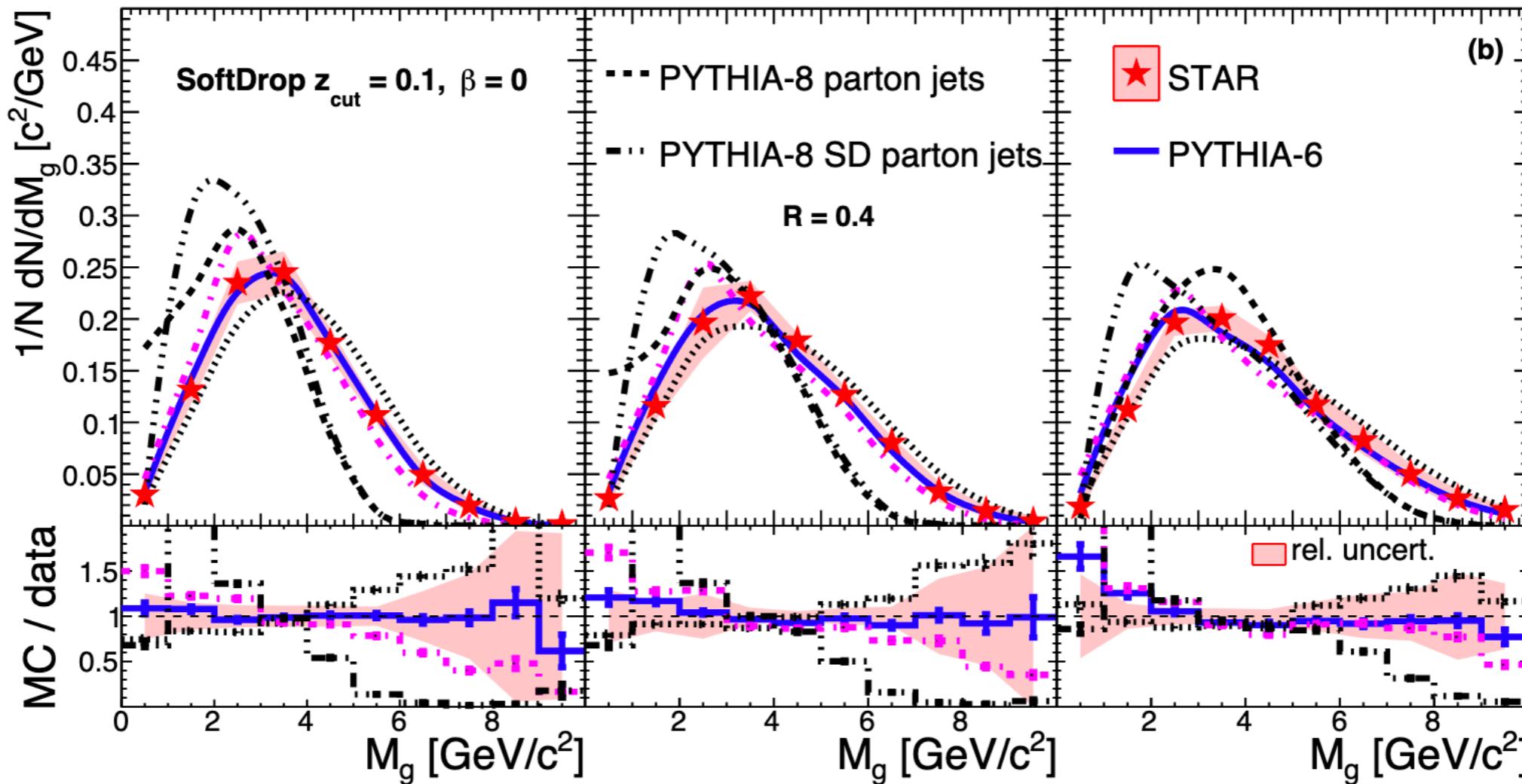
- Follow the hardest branch - Iterative SoftDrop
- Following all branches - Recursive SoftDrop

$n_{SD}$ ,  $z_g^n$ ,  $R_g^n$



# Groomed Jet Mass

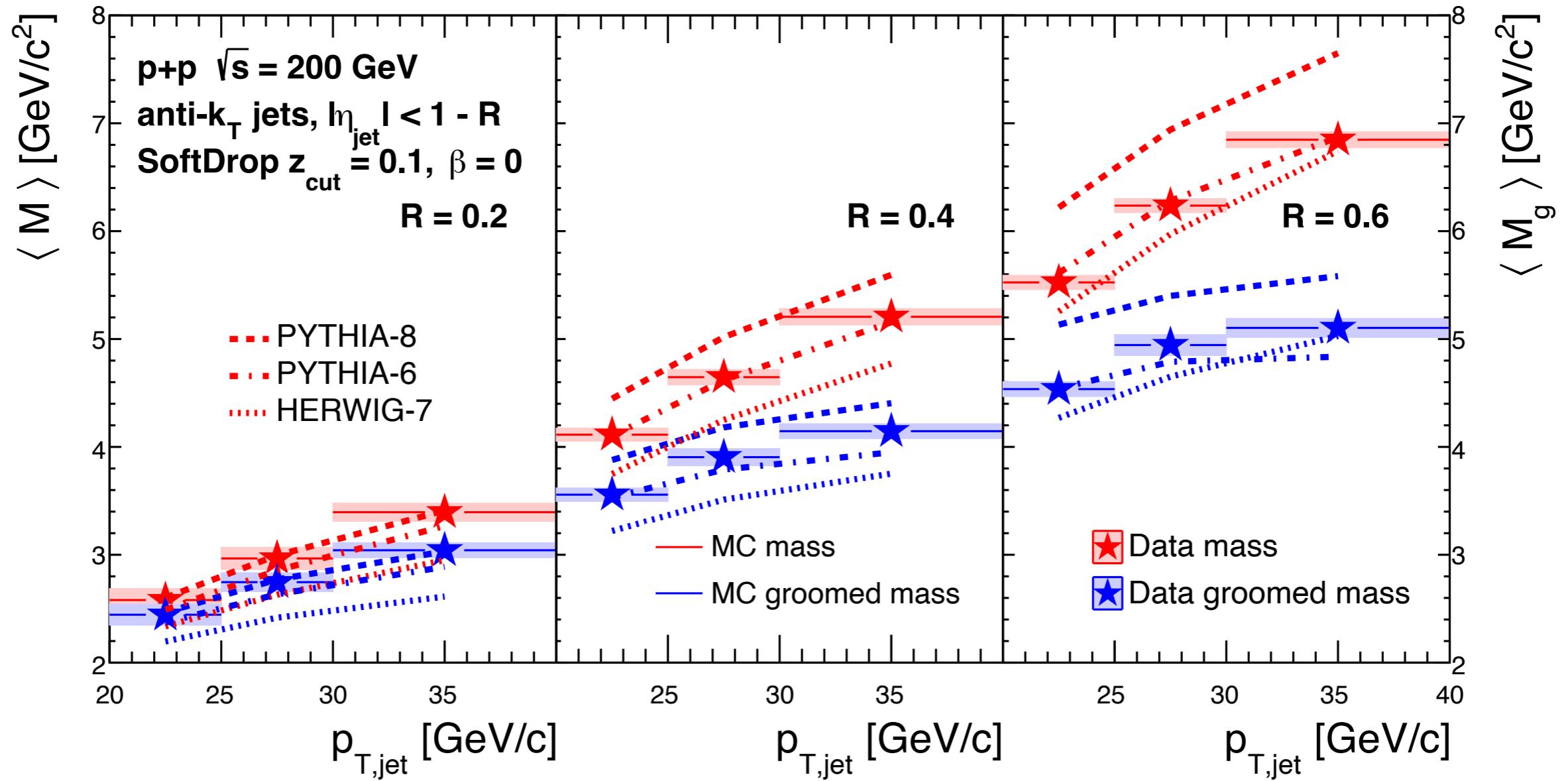
STAR arXiv:2103.13286, Accepted in PRD



RHIC-tuned **PYTHIA-6** describes **data**  
**HERWIG-7** under-predicts and **PYTHIA-8** over-predicts  
**Mass (angularity)  $\sim z\theta^2$**  Can we isolate these two scales in jets?

# Evolution of jet mass as a function of jet momenta and radii

STAR arXiv:2103.13286, Accepted in PRD



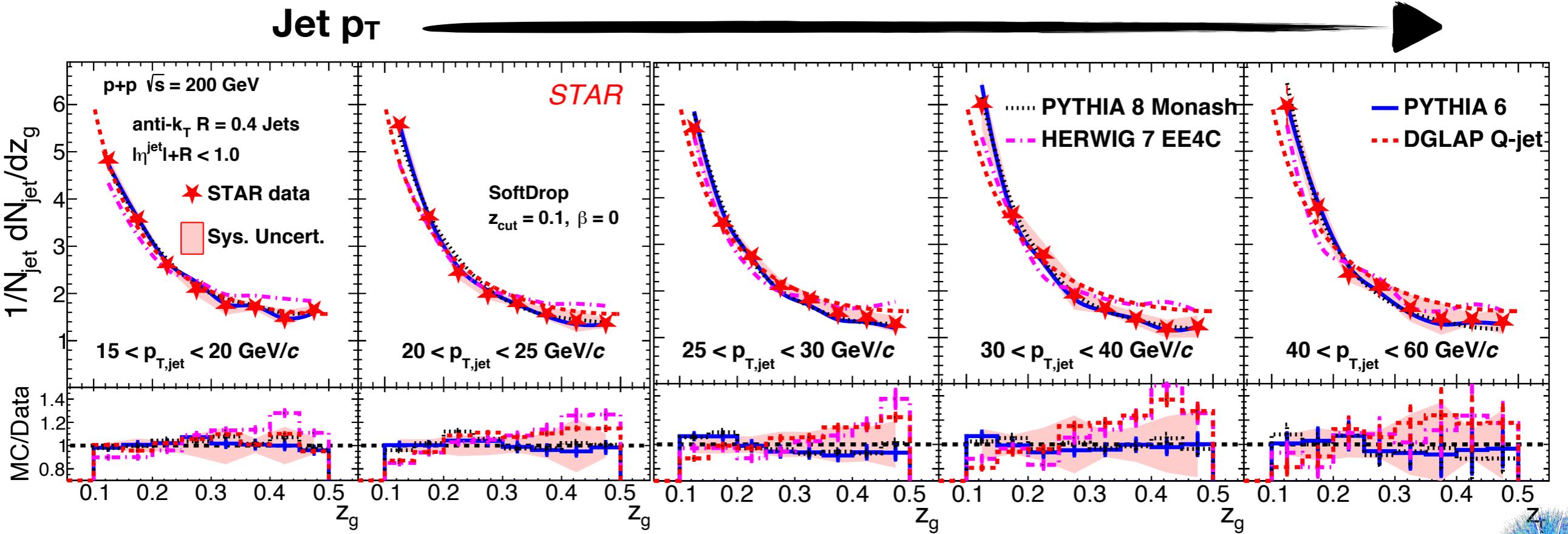
Increase in jet mass with increasing  $p_T$  and  $R$  is reduced with grooming - reduces overall impact of non-perturbative contributions to jets

# SoftDrop

$$z_g = \frac{\min(p_{T,1}, p_{T,2})}{p_{T,1} + p_{T,2}}$$

- ★ Recover the universal  $1/z$  behavior starting from  $p_T \sim 25 \text{ GeV}/c$
- ★ PYTHIA-6 and PYTHIA-8 describe data
- ★ HERWIG-7 predicts harder splitting

STAR, Phys.Lett.B 811 (2020)

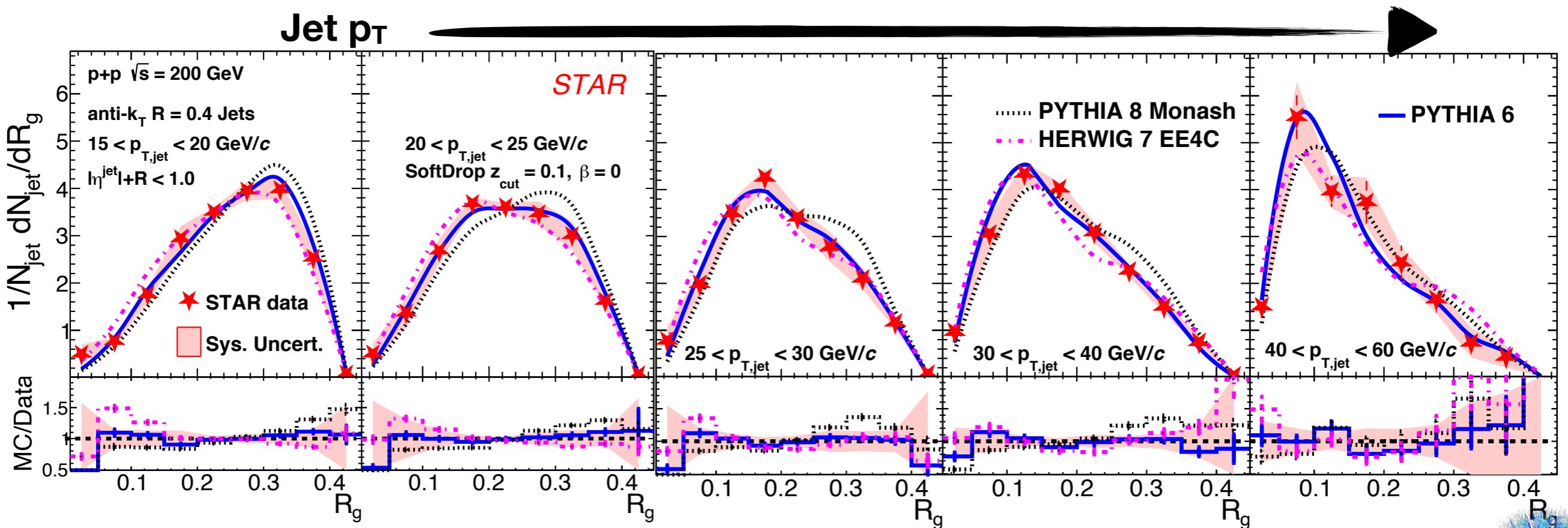


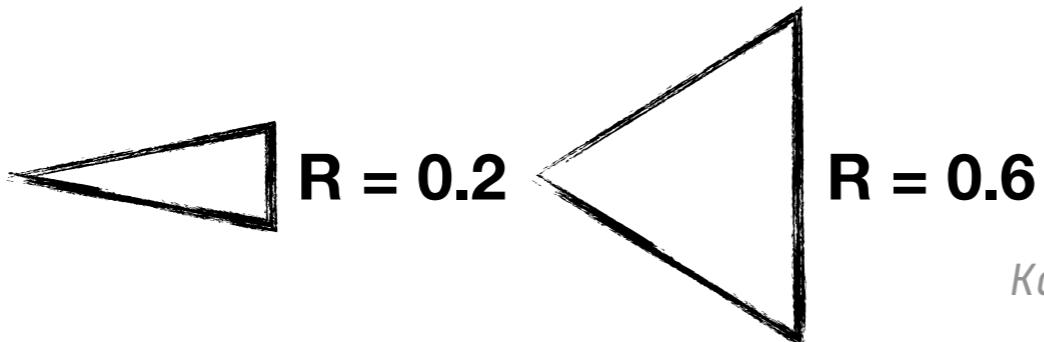
# SoftDrop $R_g$

$R_g = \Delta R(1,2)$

- ★  $R_g$  reflects momentum-dependent narrowing of jet structure
- ★ PYTHIA-6 describes data
- ★ PYTHIA-8 predicts larger groomed jet angular scale

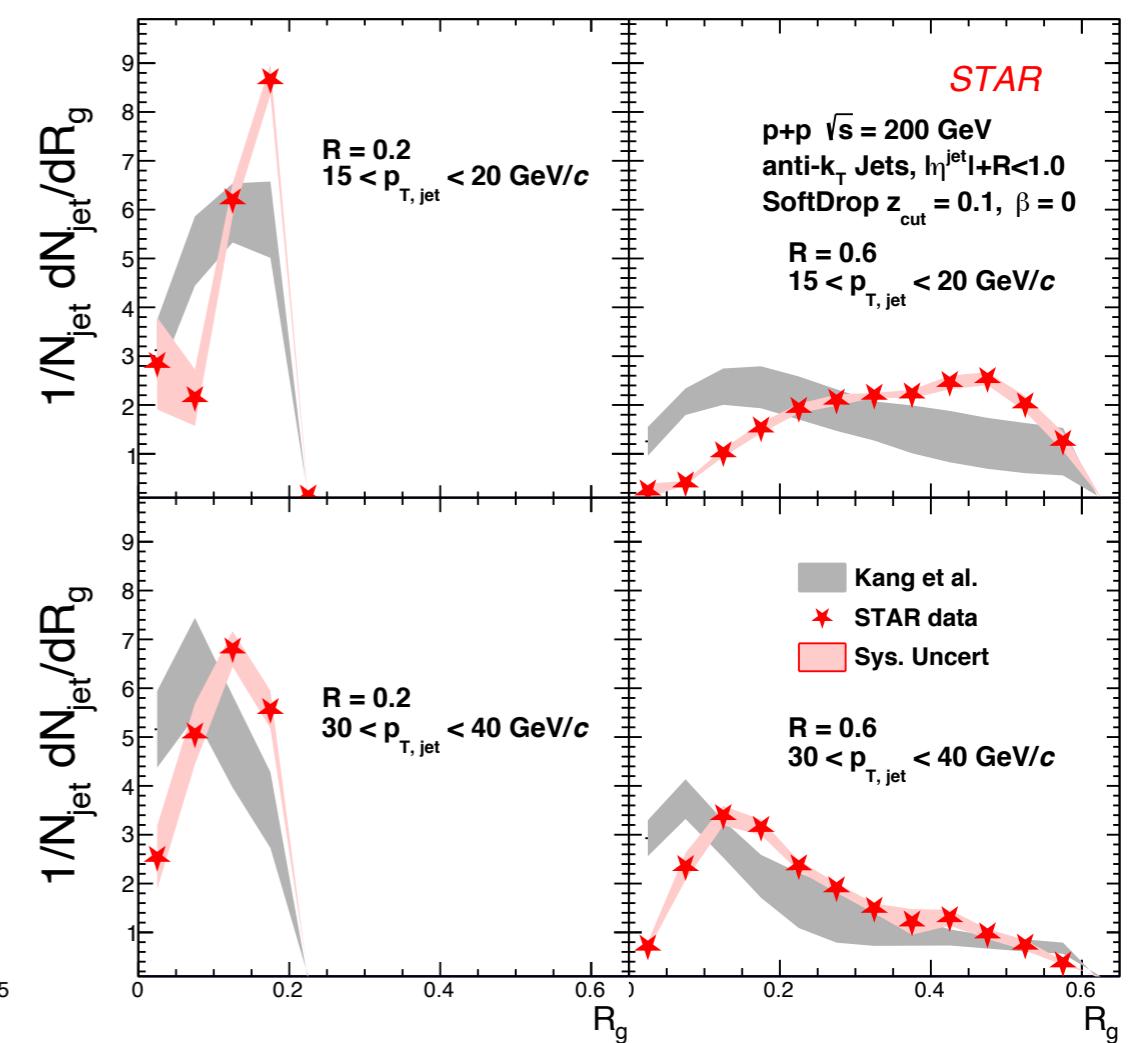
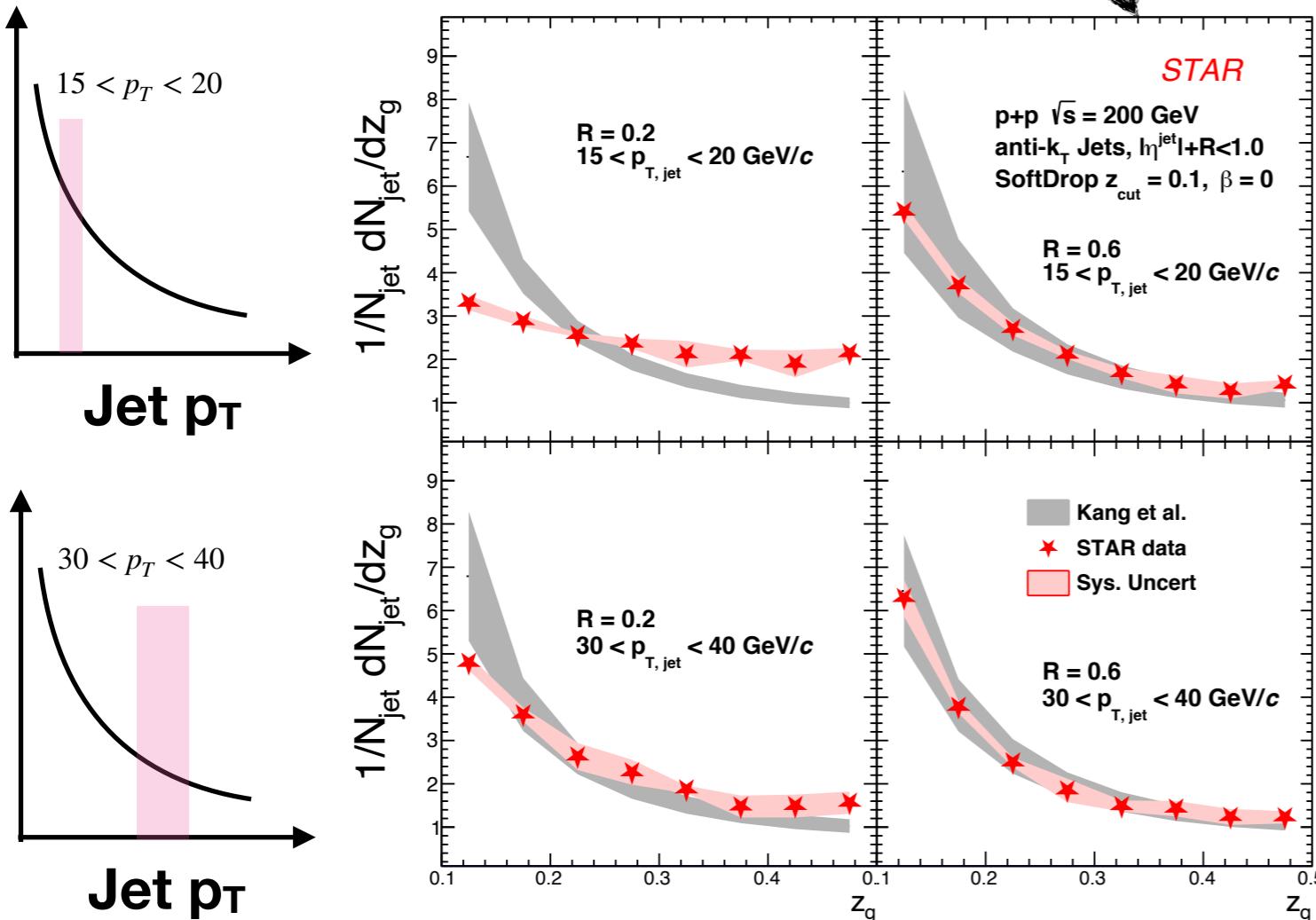
STAR, Phys.Lett.B 811 (2020)



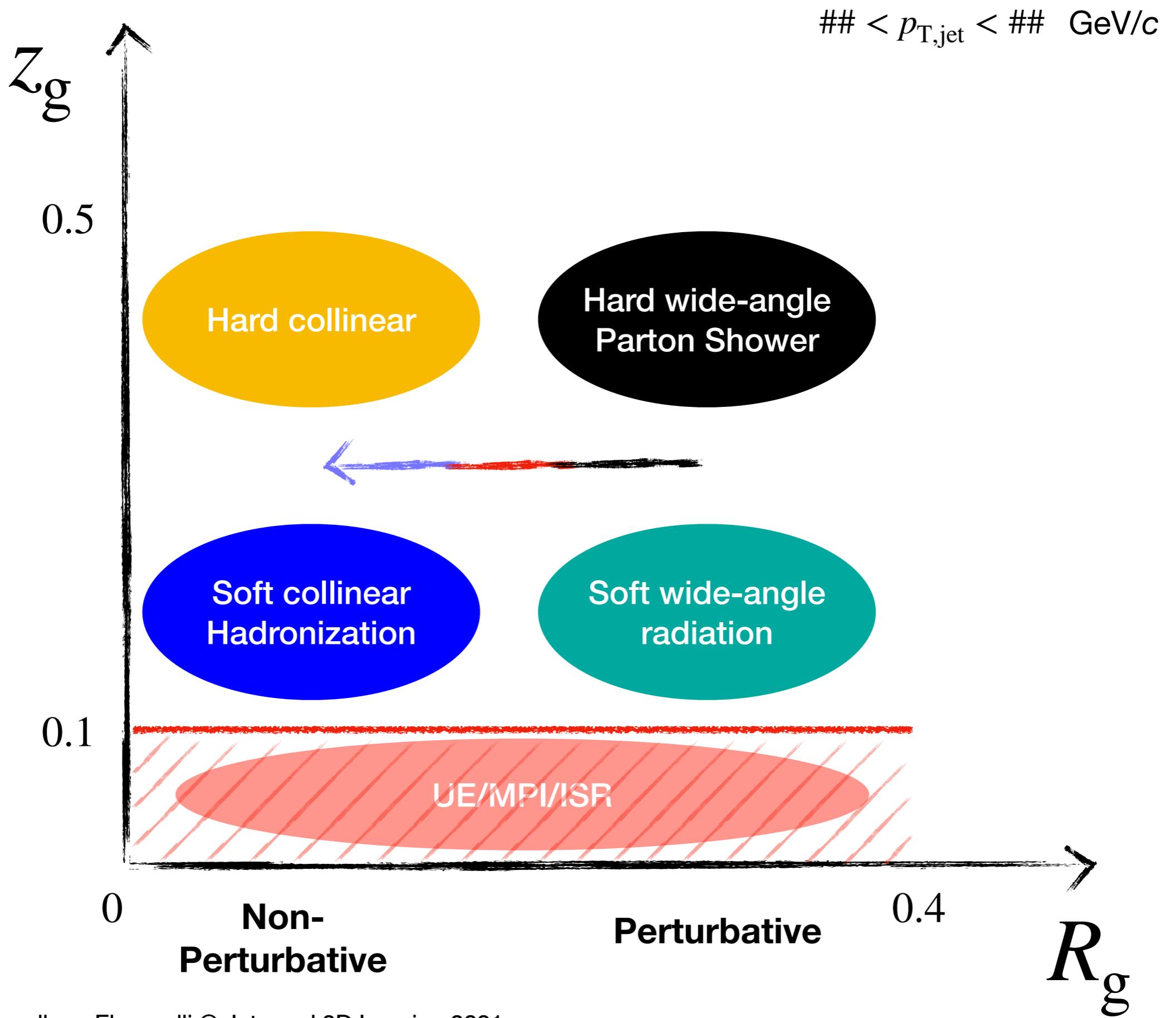


STAR, Phys.Lett.B 811 (2020)

Kang, Lee, Liu, Neill and Ringer, JHEP (2020)

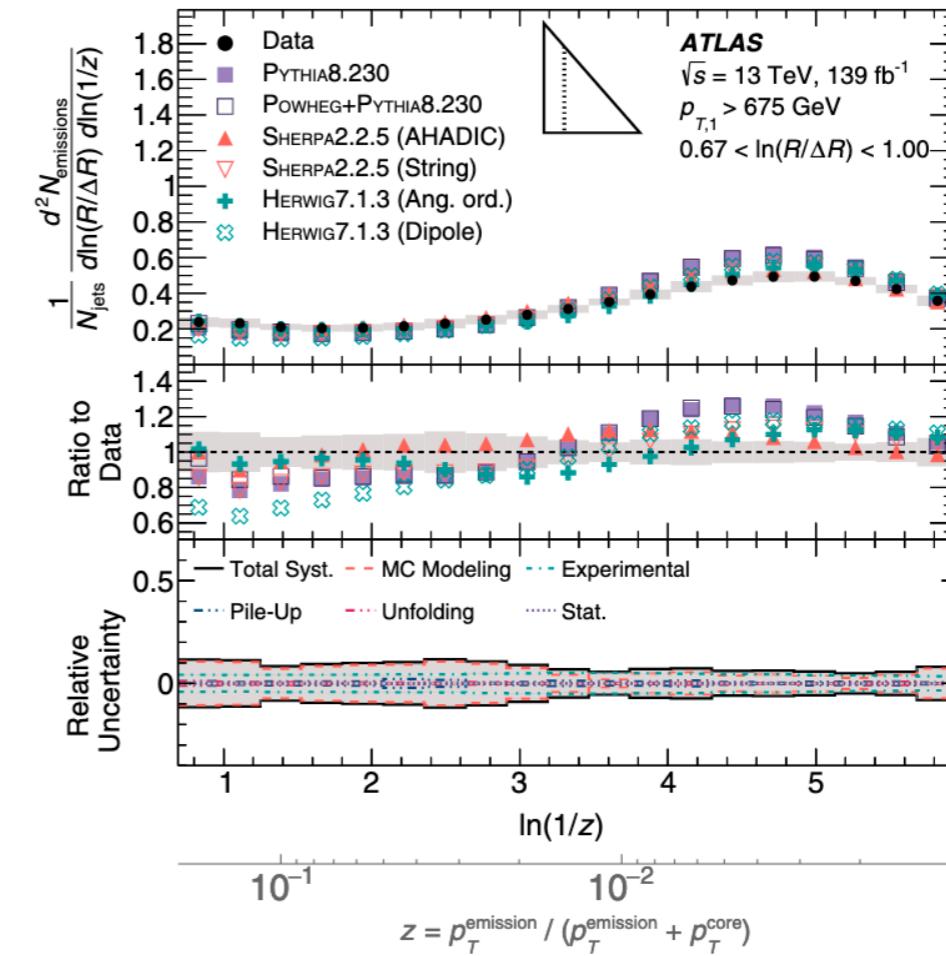
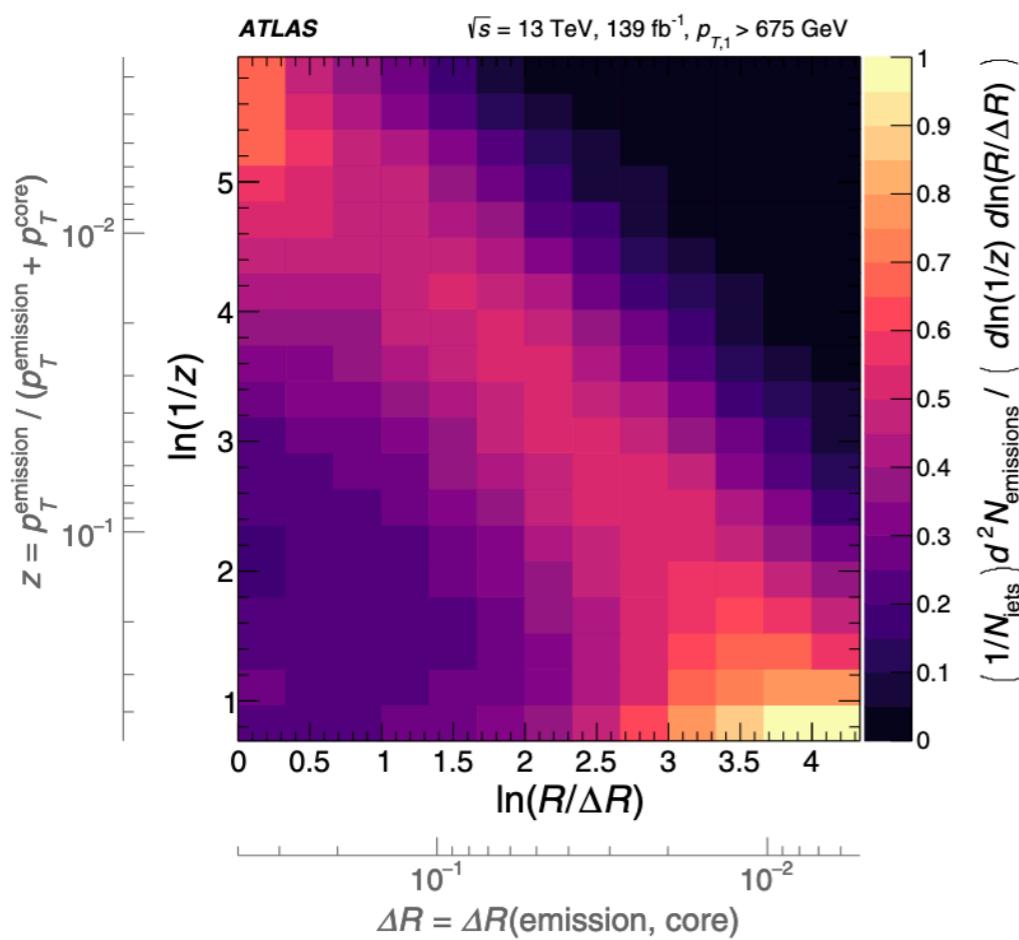


- NLL calculations (w/o non-perturbative corrections) matches data at large jet  $R$  and  $p_T$
- Significantly worse for jets of narrow  $R$  and low  $p_T$  - tighter constraints on jet splittings



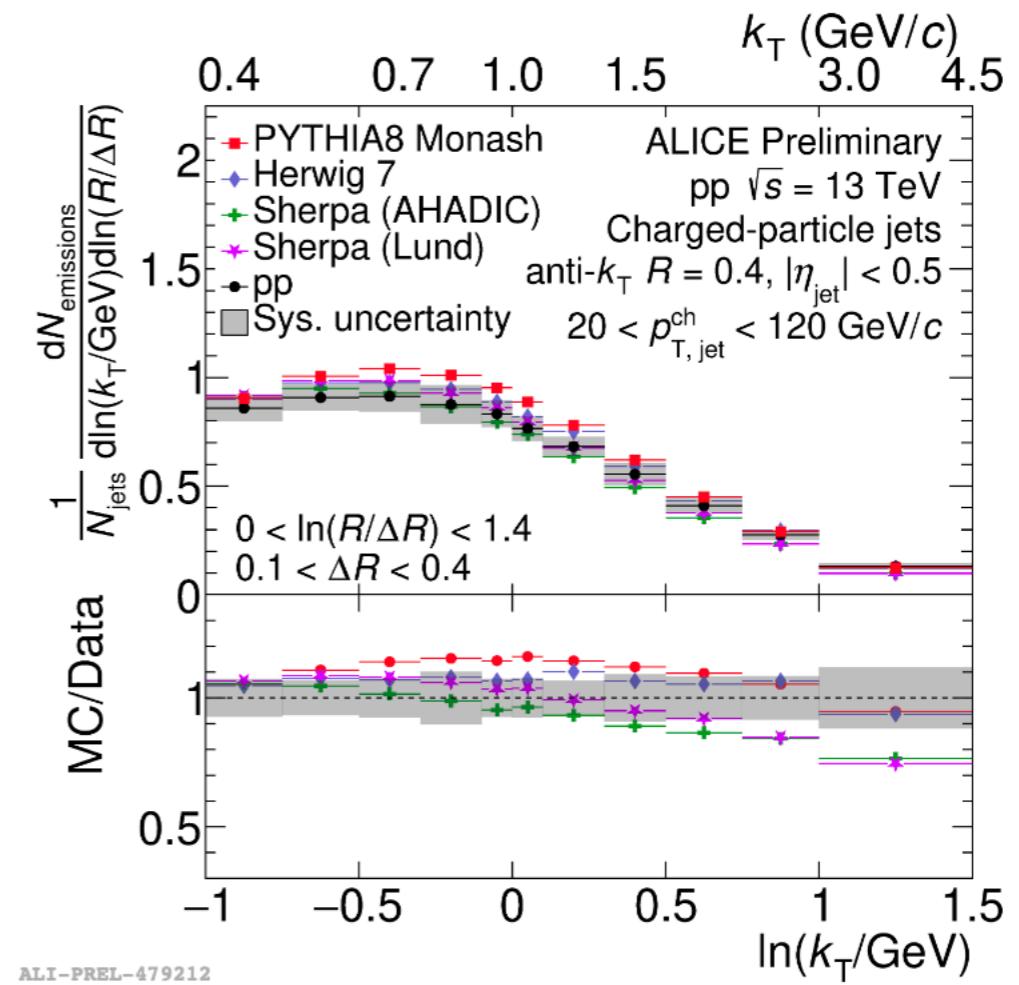
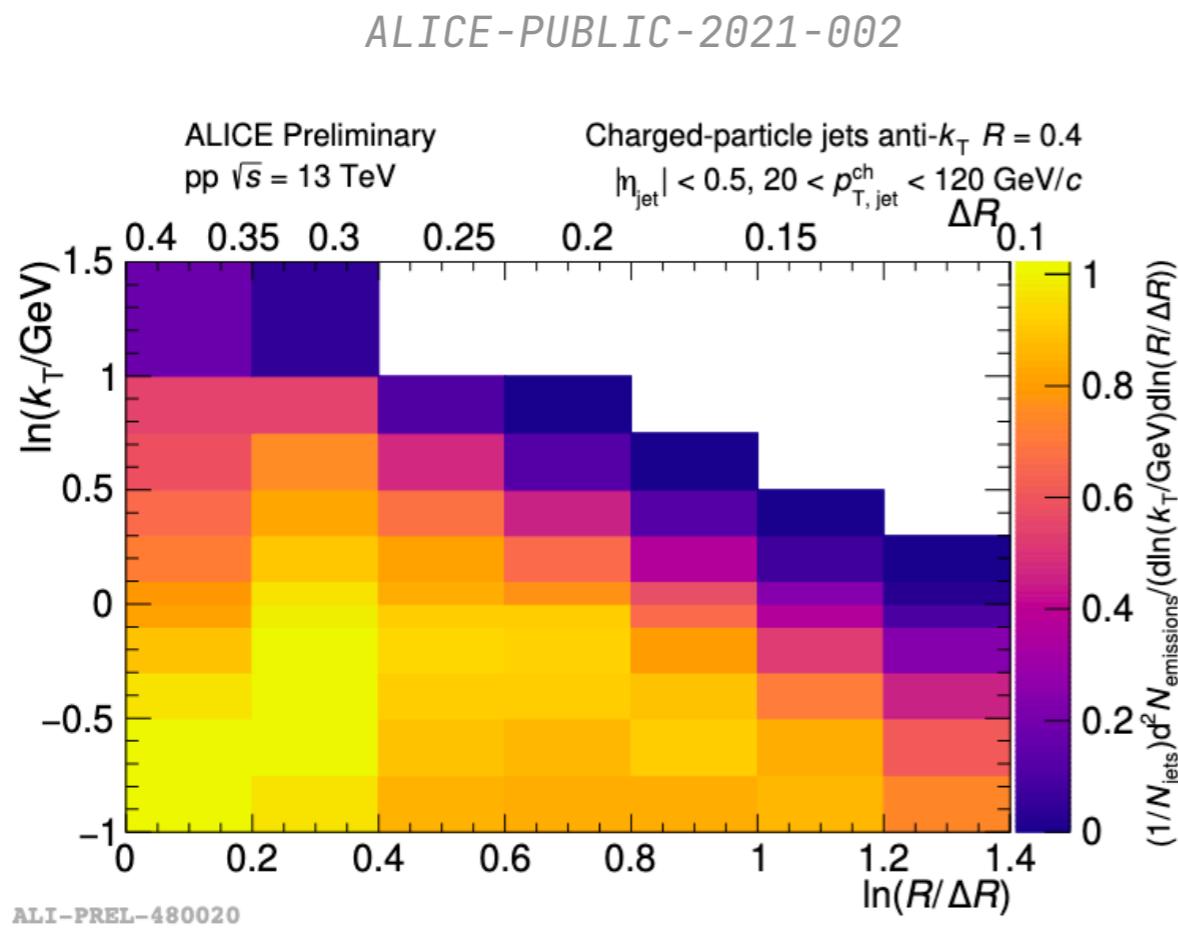
# Recent measurements of Lund Plane and their projections at the LHC

ATLAS, Phys. Rev. Lett. 124, 222002 (2020)



- Each split along the harder branch makes an entry here in the 2D Lund plane
- Comparison with particle level MC w/ varied shower/hadronization models showcase differences

# Recent measurements of Lund Plane and their projections at the LHC



- Lower  $p_T$  jets at ALICE (20 - 120 GeV) also show interesting differences for large  $k_T$  splits
- Lund plane integrates over splits - can we measure the evolution of these observables along the jet shower?

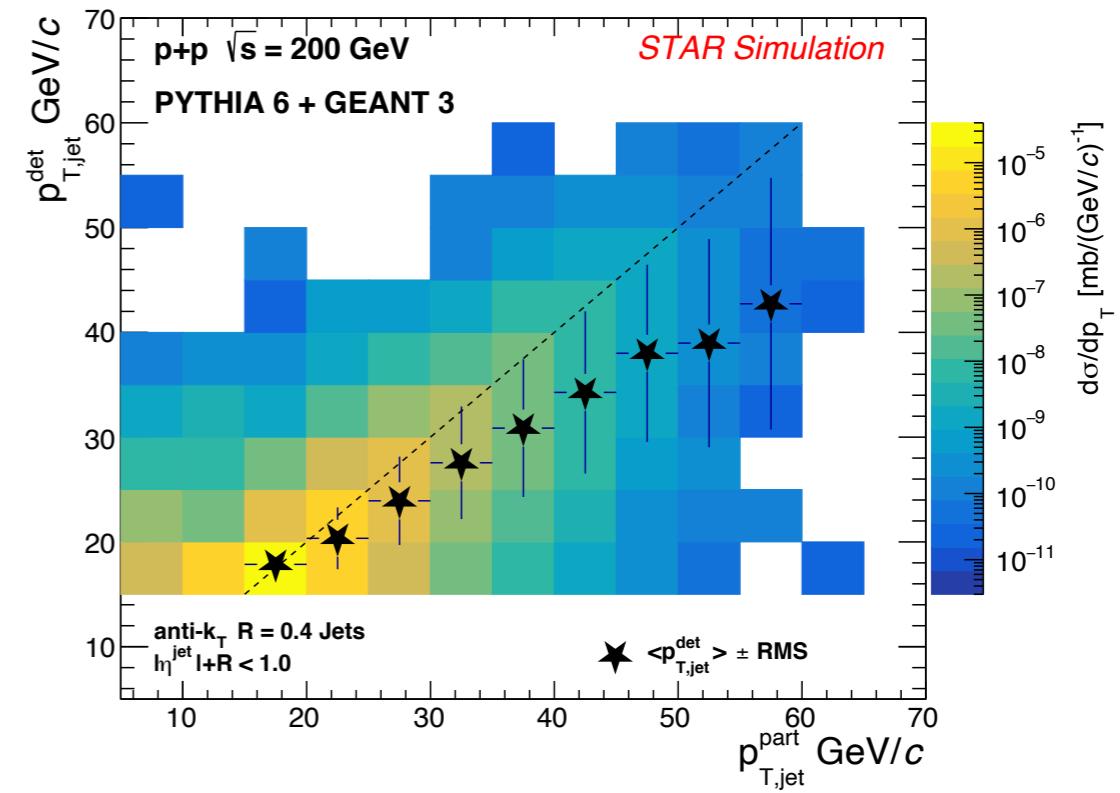
# Corrections in 3-D

## Jet $p_T, z_g, R_g$

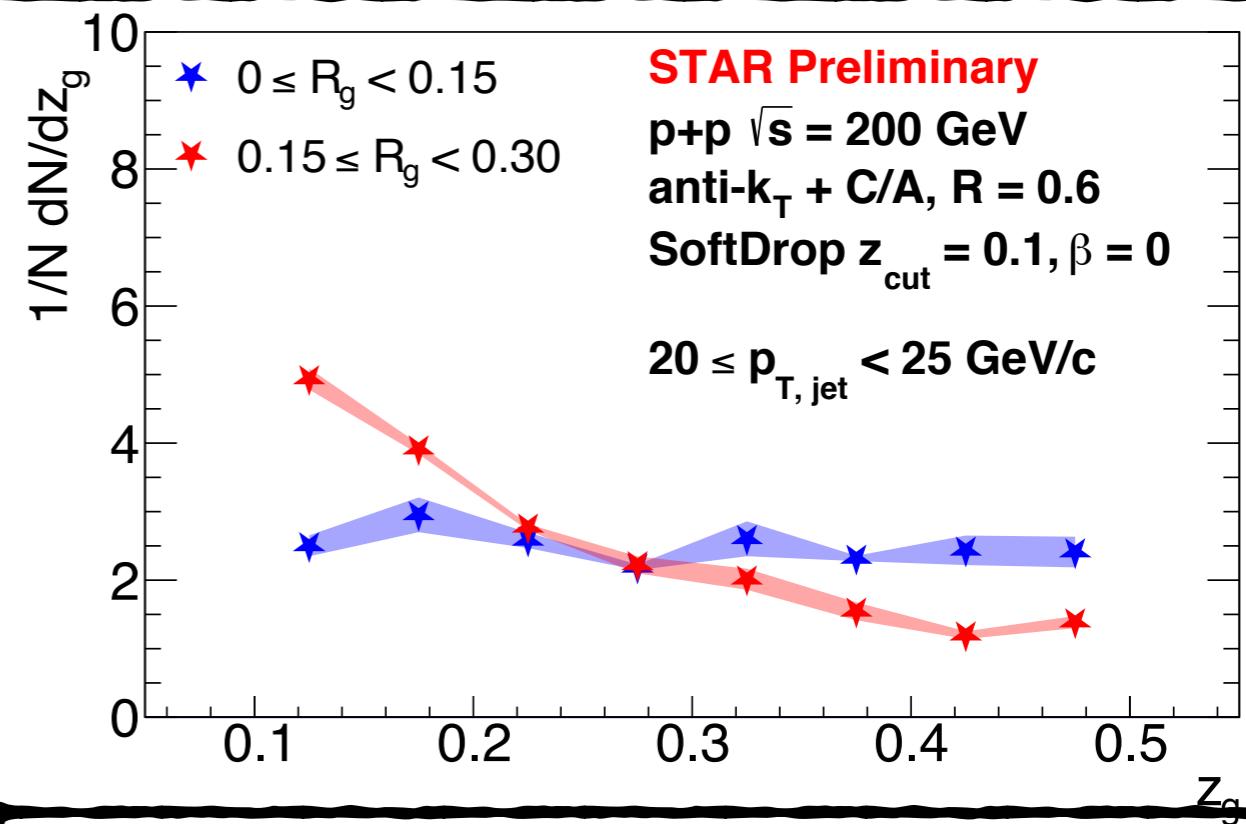
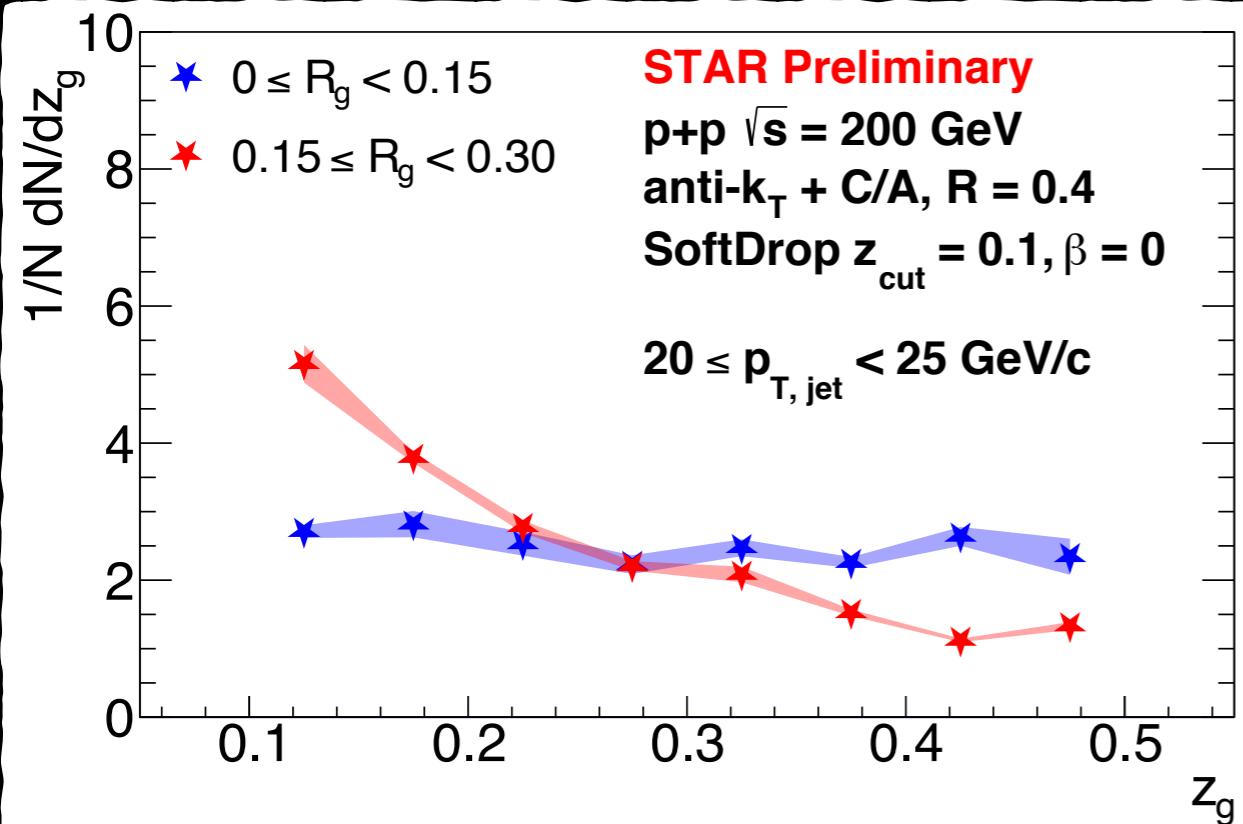
- Correlations between  $z_g$  vs.  $R_g$  at fixed detector-level jet  $p_T$   
unfolded by iterative Bayesian procedure =  $U(z_g, R_g) |_{\text{det}-p_T}$

- Since results are presented for true/particle-level jet momentum selections, corrections are done by weighted sum according to the  $p_T$  response matrix

$$\sum_{i \in \text{det}-p_T} \omega_i \cdot U(z_g, R_g) |_{\text{det}-p_T}$$

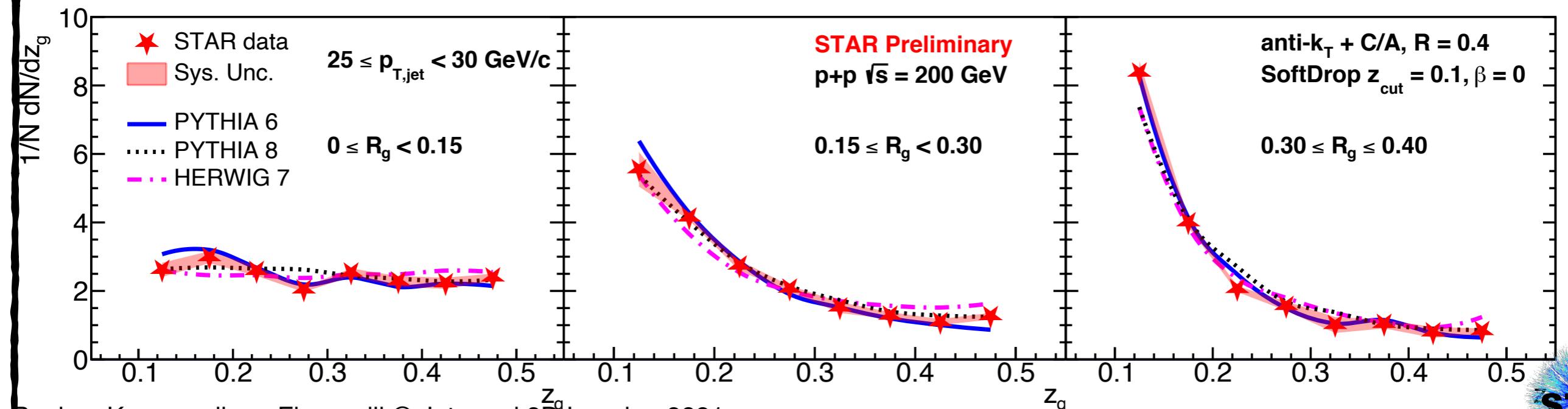


Details on systematic uncertainties available in backup



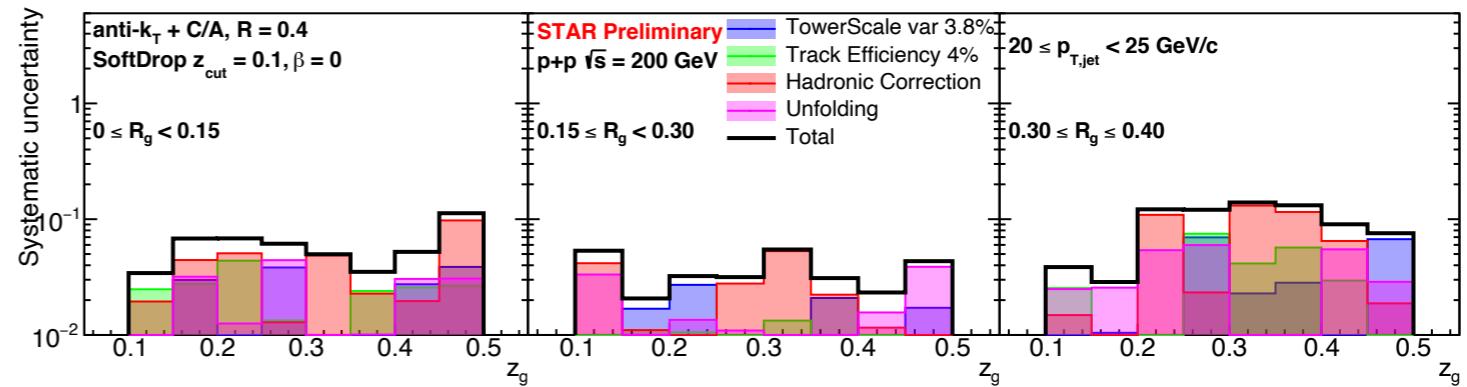
- No significant differences in substructure due to jet radius selections

- Leading order monte carlo models reproduce the evolution with different hadronization models

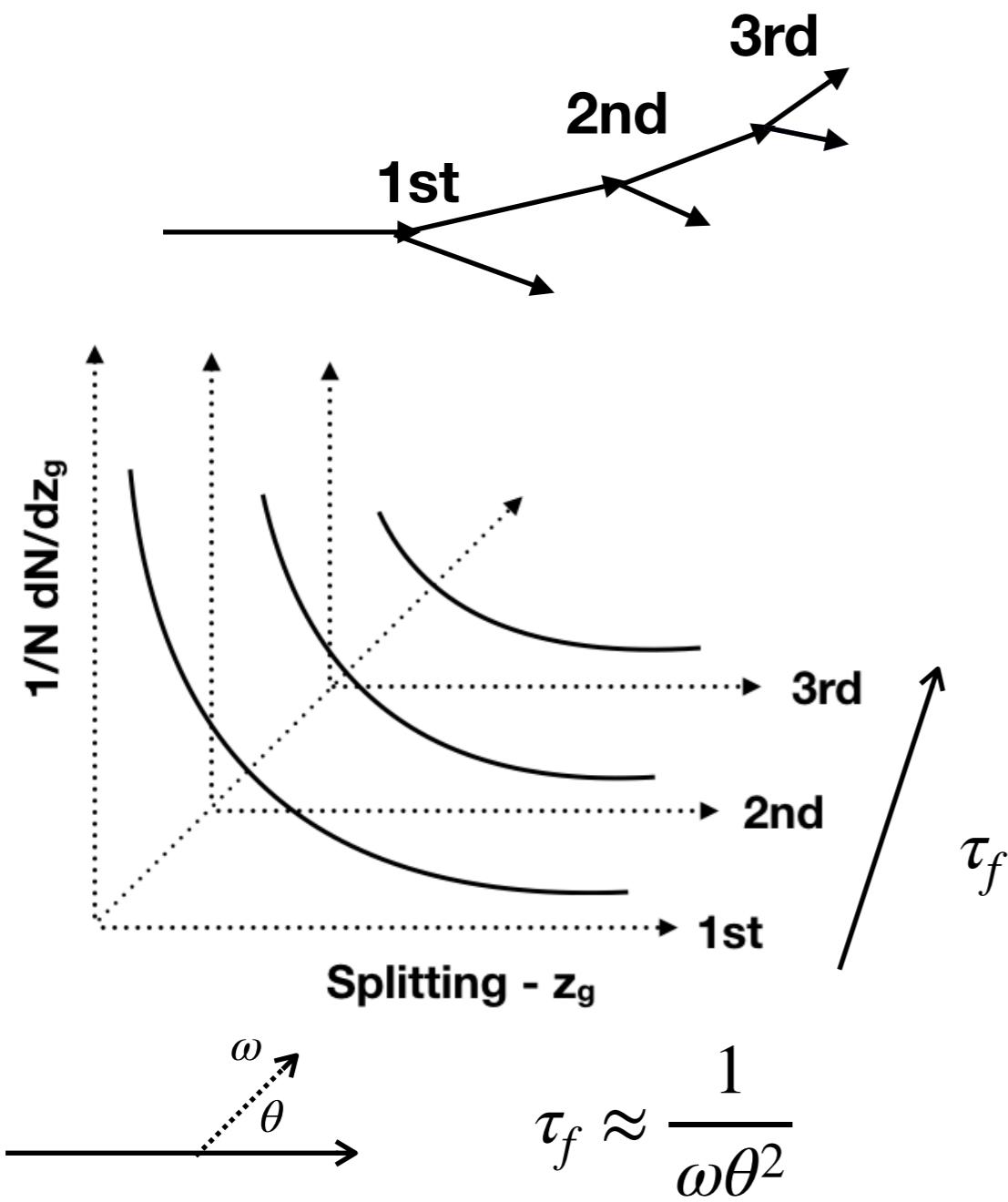


# Systematic Uncertainties

- Tracking efficiency : 4%
- Tower energy scale : 3.8%
- Hadronic correction (Matched track-tower energy subtraction) : 50% - 100%
- Bayesian unfolding iteration parameter : 2 - 6
- Prior shape variation : Priors reweighed at 1st, 2nd and 3rd split as seen in PYTHIA 6 vs PYTHIA 8 and HERWIG 7
- Split Matching criteria :  $\Delta R < 0.075, 0.1, 0.125$
- Variation in truth level shape correction for trigger and jet finding efficiencies via differences observed in PYTHIA 6 vs PYTHIA 8 and HERWIG 7



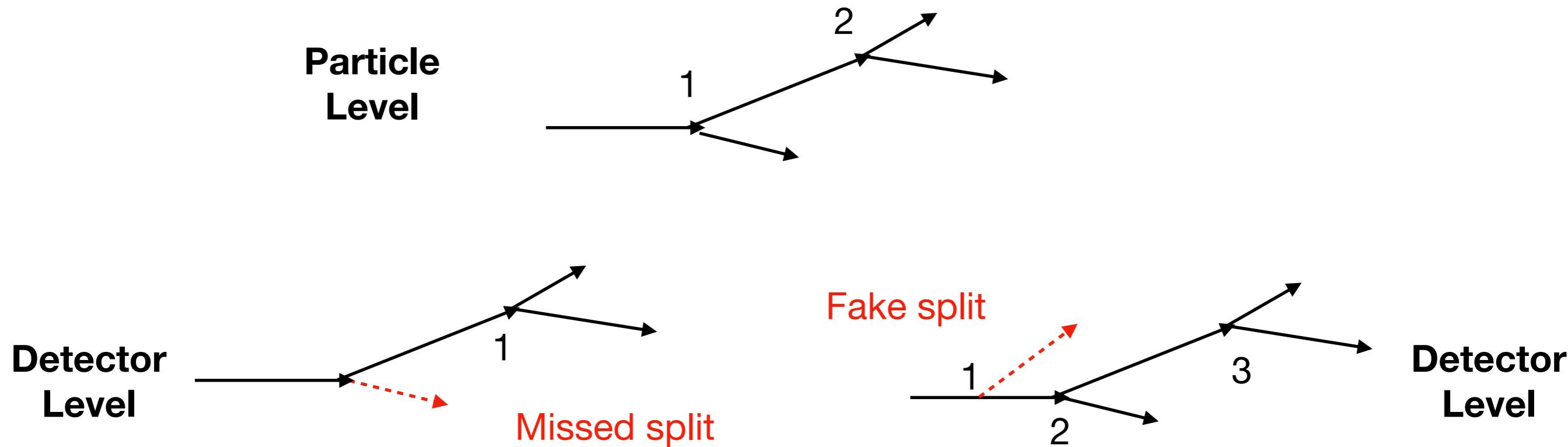
# Measure the splittings along the jet clustering tree



- Enables a study of self-similarity and effect of restricting available phase space for radiation due to virtuality evolution
- Given a jet ( $p_T^{\text{jet}}$ ) what are the  $z_g, R_g$  at 1st, 2nd and 3rd splits? **Follow a jet**
  - Compare these distributions at varying jet kinematics
  - Indirect constraint on splitting kinematics
- Given a split ( $p_T^{\text{initiator}}$ ), what are the  $z_g, R_g$  for 1st, 2nd and 3rd splits? **Follow a split**
  - Compare these at varying initiator kinematics (direct handle on splits)
  - Indirect constraint on jet kinematics

# Corrections in 3-D Jet/Initiator $p_T, z_g/R_g, n$

Finite detector efficiency and resolution can alter the splits that are reconstructed in the detector



- Observables ( $p_{T,\text{jet/initiator}}, z_g, R_g$ ) at a given split are smeared
- Splitting hierarchy also modified going from particle to detector level jets

Details of unfolding and systematic uncertainties available in backup



Jet Finding and SoftDrop

Yes

No

Discard Event

Split Matching

Yes

No

Missed Split

JP2 trigger

Pass

Fail

Trigger  
Efficiency  
(Misses)

Jet Finding and SoftDrop

Yes

No

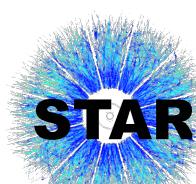
Jet Finding  
Efficiency

Split Matching

Yes

No

Fake Split



# Shape correction

Particle Level Split #

5

4

3

2

1

0

-1

-2

-2

-1

0

1

2

3

4

5

Detector Level Split #

Unmatched splits/jets  
via matching criterion

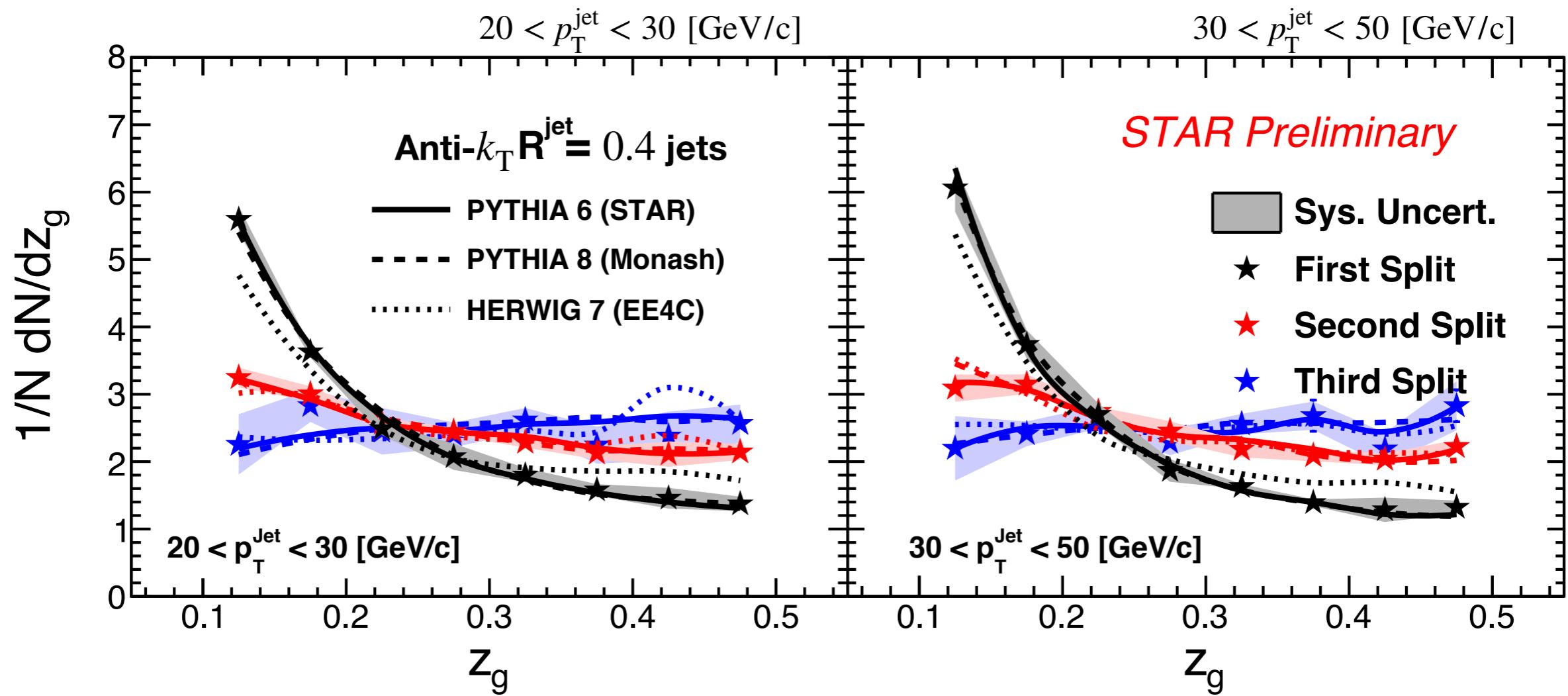
Unmatched shape

Split Matching done via  
geometric matching

Trigger Inefficiency  
no matching geant  
event for pythia event

Particle level shape  
correction (inclusive)

# Comparisons with leading order MC - $z_g$ for various jet $p_T$



- **Flattening of the splitting  $z_g$  as we increase split number** captured by the MC
- Small differences between PYTHIA and HERWIG seen in the **first** split appear to be reduced at the **second/third** splits